PAUL SCHERRER INSTITUT





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# Introduction to group theory and the classification of [superconducting] states

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# Outline

#### Brief introduction to group theory concepts:

Group  $\Rightarrow$  Conjugacy Classes  $\Rightarrow$  Group Representation  $\Rightarrow$  Character  $\Rightarrow$  Irreducible Representations

#### **Crystallographic Point Groups:**

- $\Rightarrow$  SC order parameter classification [Sigrist-Ueda]
- $\Rightarrow$  Conventional/unconventional
- $\Rightarrow$  Nematic/Chiral

#### **Beyond the Sigrist-Ueda Classification:**

- ⇒ Multiple internal DOFs (orbitals/layers/sublattices)
- $\Rightarrow$  Nonsymmorphic symmetries

#### **Introduction to Group Theory**

# Bibliography

M. Hamermesh, *Group Theory and its Application to Physical Problems*, Addison-Wesley (1962);

C. J. Bradley and A. P. Cracknell, *The Mathematical Theory of Symmetry in Solids: Representation Theory for Point Groups and Space Groups*, Claredon Press (1972);

M.S. Dresselhaus, G. Dresselhaus, and A. Jorio, *Group Theory Application to the Physics of Condensed Matter*, Springer (2008).







Group

#### What is a group?

# What is a group?

**Definition:** A group  $\mathbf{G}$  is a set of elements together with a composition law (.), also referred to as product or multiplication law, such that:

1. The product of any two elements is a member of the group:

if A and  $B \in \mathbf{G}$ , then  $A.B \in \mathbf{G}$ ;

2. The product is associative:

A.(B.C) = (A.B).C for all  $A, B, C \in \mathbf{G}$ ;

3. There exists a unique identity element E:

E.A = A.E = A for all  $A \in \mathbf{G}$ ;

4. Every element has a unique inverse:

**Example A:** The integer numbers (...-2,-1,0,1,2,...) with the operation of addition (+) is called the *additive group of the integers*. The requirements above hold:

1. The composition law (here addition) of any two elements is a member of the group:

 $1 + 1 = 2, 2 + 7 = 9, (-5) + 3 = (-2), (-1) + (-3) = (-4), \dots$ 

2. The composition is associative:

$$1 + 5 + (-3) = (1 + 5) + (-3) = 6 + (-3) = 3$$
$$1 + 5 + (-3) = 1 + [5 + (-3)] = 1 + 2 = 3$$

3. There exists a unique identity element E = 0:

 $1 + 0 = 1, 3 + 0 = 3, (-5) + 0 = (-5), \dots$ 

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Order of the group: number of elements in the group [The additive group of the integers is infinite]

**Example B:** The group of transformations of the equilateral triangle. The group is composed of the identity, rotations by  $120^{\circ}$   $(R_1)$  and  $240^{\circ}$   $(R_2)$  around the axis passing through the center of the triangle (coming out of the page), and reflections at three different mirror planes which pass though the center and the triangle's edges:  $M_i$  with i = 1, 2, 3, as indicated in Fig. 1.



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six operations in the group:  $E, R_1, R_2, M_1, M_2, M_3$ .

**Order of the group:** number of elements in the group [The group of symmetries of the equilateral triangle has order 6]

 The "product" (here the composition) of any two elements is a member of the group Convention: Apply first the right most operation.



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(you can check the remaining combinations)

1. The "product" (here the composition) of any two elements is a member of the group (note that the right most operation is the one to be applied first):



(you can check the remaining combinations)

1. The product of any two elements is a member of the group:

if A and  $B \in \mathbf{G}$ , then  $A.B \in \mathbf{G}$ ;

There is a total of 36 pairs of operations to be checked. You can check that all combinations result in one of the six operations in the group:  $E, R_1, R_2, M_1, M_2, M_3$ .



	Ε	$\mathbf{R_1}$	$\mathbf{R_2}$	$\mathbf{M_1}$	$M_2$	$M_3$
$\mathbf{E}$	E	$R_1$	$R_2$	$M_1$	$M_2$	$M_3$
$\mathbf{R_1}$	$R_1$	$R_2$	E	$M_2$	$M_3$	$M_1$
$\mathbf{R_2}$	$R_2$	E	$R_1$	$M_3$	$M_1$	$M_2$
$\mathbf{M_1}$	$M_1$	$M_3$	$M_2$	E	$R_2$	$R_1$
$\mathbf{M_2}$	$M_2$	$M_1$	$M_3$	$R_1$	E	$R_2$
$\mathbf{M_3}$	$M_3$	$M_2$	$M_1$	$R_2$	$R_1$	E

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	$\mathbf{E}$	$\mathbf{R_1}$	$\mathbf{R_2}$	$\mathbf{M_1}$	$M_2$	$M_3$
E	E	$R_1$	$R_2$	$M_1$	$M_2$	$M_3$
$\mathbf{R_1}$	$R_1$	$R_2$	E	$M_2$	$M_3$	$M_1$
$\mathbf{R_2}$	$R_2$	E	$R_1$	$M_3$	$M_1$	$M_2$
$\mathbf{M_1}$	$M_1$	$M_3$	$M_2$	E	$R_2$	$R_1$
$\mathbf{M_2}$	$M_2$	$M_1$	$M_3$	$R_1$	E	$R_2$
$\mathbf{M_3}$	$M_3$	$M_2$	$M_1$	$R_2$	$R_1$	E

#### **Multiplication Table**



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A.(B.C) = (A.B).C for all  $A, B, C \in \mathbf{G}$ ;

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	$\mathbf{E}$	$\mathbf{R_1}$	$\mathbf{R_2}$	$\mathbf{M_1}$	$\mathbf{M_2}$	$M_3$
$\mathbf{E}$	E	$R_1$	$R_2$	$M_1$	$M_2$	$M_3$
$\mathbf{R}_1$	$R_1$	$R_2$	E	$M_2$	$M_3$	$M_1$
$\mathbf{R_2}$	$R_2$	E	$R_1$	$M_3$	$M_1$	$M_2$
$\mathbf{M_1}$	$M_1$	$M_3$	$M_2$	E	$R_2$	$R_1$
$M_2$	$M_2$	$M_1$	$M_3$	$R_1$	E	$R_2$
$M_3$	$M_3$	$M_2$	$M_1$	$R_2$	$R_1$	E

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Homework!

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Ε	E	$R_1$	$R_2$	$M_1$	$M_2$	$M_3$
$\mathbf{R_1}$	$R_1$	$R_2$	E	$M_2$	$M_3$	$M_1$
$\mathbf{R_2}$	$R_2$	E	$R_1$	$M_3$	$M_1$	$M_2$
$\mathbf{M_1}$	$M_1$	$M_3$	$M_2$	E	$R_2$	$R_1$
$M_2$	$M_2$	$M_1$	$M_3$	$R_1$	E	$R_2$
$M_3$	$M_3$	$M_2$	$M_1$	$R_2$	$R_1$	E

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	E	$\mathbf{R_1}$	$\mathbf{R_2}$	$\mathbf{M_1}$	$\mathbf{M_2}$	$\mathbf{M_3}$
$\mathbf{E}$	E	$R_1$	$R_2$	$M_1$	$M_2$	$M_3$
$\mathbf{R}_1$	$R_1$	$R_2$	E	$M_2$	$M_3$	$M_1$
$\mathbf{R_2}$	$R_2$	E	$R_1$	$M_3$	$M_1$	$M_2$
$\mathbf{M_1}$	$M_1$	$M_3$	$M_2$	E	$R_2$	$R_1$
$M_2$	$M_2$	$M_1$	$M_3$	$R_1$	E	$R_2$
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$\mathbf{R_1}$	$R_1$	$R_2$	E	$M_2$	$M_3$	$M_1$
$\mathbf{R_2}$	$R_2$	E	$R_1$	$M_3$	$M_1$	$M_2$
$\mathbf{M_1}$	$M_1$	$M_3$	$M_2$	E	$R_2$	$R_1$
$\mathbf{M_2}$	$M_2$	$M_1$	$M_3$	$R_1$	E	$R_2$
$\mathbf{M_3}$	$M_3$	$M_2$	$M_1$	$R_2$	$R_1$	E

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given  $A \in \mathbf{G}$ , there exists an element  $A^{-1}$  such that  $A \cdot A^{-1} = A^{-1} \cdot A = E$ .

	E	$\mathbf{R_1}$	$\mathbf{R_2}$	$\mathbf{M_1}$	$M_2$	$\mathbf{M_3}$
E	E	$R_1$	$R_2$	$M_1$	$M_2$	$M_3$
$\mathbf{R_1}$	$R_1$	$R_2$	E	$M_2$	$M_3$	$M_1$
$\mathbf{R_2}$	$R_2$	E	$R_1$	$M_3$	$M_1$	$M_2$
$\mathbf{M_1}$	$M_1$	$M_3$	$M_2$	E	$R_2$	$R_1$
$\mathbf{M_2}$	$M_2$	$M_1$	$M_3$	$R_1$	E	$R_2$
$\mathbf{M_3}$	$M_3$	$M_2$	$M_1$	$R_2$	$R_1$	E

#### C<sub>3v</sub> point group [isomorphic to S<sub>3</sub>]

**Conjugate Elements:** Two elements  $G_1$  and  $G_2$  are said to be conjugate if there exists an element G in **G** such that  $G_1 = GG_2G^{-1}$ ;



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#### Group of Symmetries of the Equilateral triangle

	$\mathbf{E}$	$\mathbf{R_1}$	$\mathbf{R_2}$	$\mathbf{M_1}$	$M_2$	$\mathbf{M_3}$
Ε	E	$R_1$	$R_2$	$M_1$	$M_2$	$M_3$
$\mathbf{R_1}$	$R_1$	$R_2$	E	$M_2$	$M_3$	$M_1$
$\mathbf{R_2}$	$R_2$	E	$R_1$	$M_3$	$M_1$	$M_2$
$M_1$	$M_1$	$M_3$	$M_2$	E	$R_2$	$R_1$
$M_2$	$M_2$	$M_1$	$M_3$	$R_1$	E	$R_2$
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**I) Identity:**  $G.E.G^{-1} = G.G^{-1}.E = E$ 

E is not conjugate to any other element

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$\mathbf{E}$	E	$R_1$	$R_2$	$M_1$	$M_2$	$M_3$
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$\mathbf{R_2}$	$R_2$	E	$R_1$	$M_3$	$M_1$	$M_2$
$M_1$	$M_1$	$M_3$	$M_2$	E	$R_2$	$R_1$
$M_2$	$M_2$	$M_1$	$M_3$	$R_1$	E	$R_2$
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**II)** Rotations:  $M_i.R_1.M_i^{-1} = M_i.R_1M_i = R_2$ 

Rotations are conjugate to each other

	$\mathbf{E}$	$\mathbf{R_1}$	$\mathbf{R_2}$	$\mathbf{M_1}$	$\mathbf{M_2}$	$\mathbf{M_3}$
$\mathbf{E}$	E	$R_1$	$R_2$	$M_1$	$M_2$	$M_3$
$\mathbf{R}_1$	$R_1$	$R_2$	E	$M_2$	$M_3$	$M_1$
$\mathbf{R_2}$	$R_2$	E	$R_1$	$M_3$	$M_1$	$M_2$
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**II)** Rotations:  $M_i.R_1.M_i^{-1} = M_i.R_1M_i = R_2$ 

Rotations are conjugate to each other

**III) Reflections:**  $R_1 \cdot M_a \cdot R_1^{-1} = R_1 \cdot M_a \cdot R_2 = M_b$  $a = \{1, 2, 3\}$  and  $b = \{3, 1, 2\}$ 

Reflections are conjugate among themselves

	$\mathbf{E}$	$\mathbf{R_1}$	$\mathbf{R_2}$	$\mathbf{M_1}$	$\mathbf{M_2}$	$\mathbf{M_3}$
$\mathbf{E}$	E	$R_1$	$R_2$	$M_1$	$M_2$	$M_3$
$\mathbf{R}_1$	$R_1$	$R_2$	E	$M_2$	$M_3$	$M_1$
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**Conjugacy classes:** The elements of a group can be split into conjugacy classes  $C_1, C_2, C_3, \ldots$  such that the following properties hold:

- 1. Every element of **G** is in some class and no element of **G** is in more than one class such that  $\mathbf{G} = C_1 + C_2 + C_3 + \dots$ ;
- 2. All elements in a given class are mutually conjugate and consequently have the same order;
- 3. An element that commutes with all other elements of the group is on a class by itself;
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**Order of an element:** the number of times the element needs to be applied to be equal to the identity.

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- 3. An element that commutes with all other elements of the group is on a class by itself;
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- An element that commutes with all other elements of the group is on a class by itself; Always the case for the identity!
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- An element that commutes with all other elements of the group is on a class by itself; Always the case for the identity!
- 4. The number of elements in a class is a divisor of the order of the group;
- $C_1 = \{E\}$  $C_2 = \{R_1, R_2\}$  $C_3 = \{M_1, M_2, M_3\}$

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**Definition:** A representation of a group **G** is a mapping D of the elements of **G** onto a set of linear operators (or matrices) with the following properties: (i) D(E) = 1, where 1 is the identity operator in the space on which the linear operator acts.

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	Ε	$\mathbf{R_1}$	$\mathbf{R_2}$	$\mathbf{M_1}$	$\mathbf{M_2}$	$\mathbf{M_3}$
$\mathbf{E}$	E	$R_1$	$R_2$	$M_1$	$M_2$	$M_3$
$\mathbf{R_1}$	$R_1$	$R_2$	E	$M_2$	$M_3$	$M_1$
$\mathbf{R_2}$	$R_2$	E	$R_1$	$M_3$	$M_1$	$M_2$
$\mathbf{M_1}$	$M_1$	$M_3$	$M_2$	E	$R_2$	$R_1$
$\mathbf{M_2}$	$M_2$	$M_1$	$M_3$	$R_1$	E	$R_2$
$\mathbf{M_3}$	$M_3$	$M_2$	$M_1$	$R_2$	$R_1$	E





Ok, but what about nontrivial representations?



#### Group of Symmetries of the Equilateral triangle

Thinking of transformations acting on the coordinates (x,y,z):





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$$R_1 = \begin{pmatrix} -1/2 & +\sqrt{3}/2 & 0\\ -\sqrt{3}/2 & -1/2 & 0\\ 0 & 0 & 1 \end{pmatrix}$$

$$M_1 = \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

	$\mathbf{E}$	$\mathbf{R_1}$	$\mathbf{R_2}$	$\mathbf{M_1}$	$M_2$	$M_3$
E	E	$R_1$	$R_2$	$M_1$	$M_2$	$M_3$
$\mathbf{R_1}$	$R_1$	$R_2$	E	$M_2$	$M_3$	$M_1$
$\mathbf{R_2}$	$R_2$	E	$R_1$	$M_3$	$M_1$	$M_2$
$\mathbf{M_1}$	$M_1$	$M_3$	$M_2$	E	$R_2$	$R_1$
$\mathbf{M_2}$	$M_2$	$M_1$	$M_3$	$R_1$	E	$R_2$
$\mathbf{M_3}$	$M_3$	$M_2$	$M_1$	$R_2$	$R_1$	E

You can check if matrices reproduce the structure of the group





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E	E	$R_1$	$R_2$	$M_1$	$M_2$	$M_3$
$\mathbf{R_1}$	$R_1$	$R_2$	E	$M_2$	$M_3$	$M_1$
$\mathbf{R_2}$	$R_2$	E	$R_1$	$M_3$	$M_1$	$M_2$
$\mathbf{M_1}$	$M_1$	$M_3$	$M_2$	E	$R_2$	$R_1$
$\mathbf{M_2}$	$M_2$	$M_1$	$M_3$	$R_1$	E	$R_2$
$\mathbf{M_3}$	$M_3$	$M_2$	$M_1$	$R_2$	$R_1$	E

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Dimension of the representation: the dimension of the space on which it acts





Thinking of transformations acting on the coordinates (x,y,z):



	E	$\mathbf{R_1}$	$\mathbf{R_2}$	$\mathbf{M_1}$	$M_2$	$M_3$
E	E	$R_1$	$R_2$	$M_1$	$M_2$	$M_3$
$\mathbf{R_1}$	$R_1$	$R_2$	E	$M_2$	$M_3$	$M_1$
$\mathbf{R_2}$	$R_2$	E	$R_1$	$M_3$	$M_1$	$M_2$
$\mathbf{M_1}$	$M_1$	$M_3$	$M_2$	E	$R_2$	$R_1$
$\mathbf{M_2}$	$M_2$	$M_1$	$M_3$	$R_1$	E	$R_2$
$\mathbf{M_3}$	$M_3$	$M_2$	$M_1$	$R_2$	$R_1$	E

You can check if matrices reproduce the structure of the group

Dimension of the representation: the dimension of the space on which it acts

**Generators of the group:** the minimal set of operations out of which the entire group can be derived [not unique]





#### Group of Symmetries of the Equilateral triangle

Thinking of transformations acting on the coordinates (x,y,z):

$$R_{1} = \begin{pmatrix} -1/2 & +\sqrt{3}/2 & 0 \\ -\sqrt{3}/2 & -1/2 & 0 \\ \hline 0 & 0 & 1 \end{pmatrix} \qquad M_{1} = \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ \hline 0 & 0 & 1 \end{pmatrix}$$

Note: The z-component never mix with the x- and y-components. This means we can divide the space in {x,y} and {z} and treat them independently. In this case we say the representation is **reducible**.



#### Group of Symmetries of the Equilateral triangle

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$$D_1(R_1) = \begin{pmatrix} -1/2 & +\sqrt{3}/2 \\ -\sqrt{3}/2 & -1/2 \end{pmatrix}$$
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**One-dimensional irreducible representation** 

$$D_2(R_1) = 1$$
  
 $D_2(M_1) = 1$ 

[Trivial representation]



Group of Symmetries of the Equilateral triangle

Two-dimensional irreducible representation

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**One-dimensional [trivial] irreducible representation** 

$$D_2(R_1) = 1$$
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Question: How can we know that we have identified all the representations?

**Character:** The characters of a group representation D are the traces of the respective linear operators (matrices)  $\chi_D(G_i) = \text{Tr}D(G_i)$ . The trace of a matrix is the sum of its diagonal elements.

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#### **Question: Have identified all the representations?**

- The number of irreducible representations, r, is equal to the number of conjugacy classes;
- The order of the group  $\mathbf{G}$ ,  $|\mathbf{G}|$ , is equal to the sum of the squares of the dimensions of the irreducible representations  $d_i$ ,  $|\mathbf{G}| = \sum_{i=1}^r d_i^2$ ;
- The characters are orthonormal:  $\sum_{i=1}^{r} n_i \chi_D^*(G_i) \chi_{D'}(G_i) = |\mathbf{G}| \delta^{DD'}$ , where  $n_i$  is the number of elements in the conjugacy class represented by  $G_i$ .

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Note that these properties can in principle be derived directly from the group structure, without thinking about any geometric realisation of the transformations!

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These are can be found in

- Bradley and Cracknell
- Bilbao crystallographic server

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bilbao crystallographic server

 $\textbf{Group} \Rightarrow \textbf{Conjugacy Classes} \Rightarrow \textbf{Group Representation} \Rightarrow \textbf{Character} \Rightarrow \textbf{Irreducible Representations}$ 

## **Crystallographic Groups**

# SC and other ordered phases emerge in...







E, R[60°], R[120°], R[180°], R[240°], R[300°]

 $=C_6 = C_3 = C_2 = C_3^{-1} = C_6^{-1}$ 



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 $=C_6 = C_3 = C_2 = C_3^{-1} = C_6^{-1}$ 










E, R[90°], R[180°], R[270°] =C<sub>4</sub> =C<sub>2</sub>



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### Character table and irreducible representations (Irrep)



$E$ $2C_4(z)$	$C_2(z)$	$2C_2(x)$	$2C_2(d)$
---------------	----------	-----------	-----------



### Character table and irreducible representations (Irrep)

TopBottom

Irrep	E	$2C_4(z)$	$C_2(z)$	$2C_2(x)$	$2C_2(d)$
$A_1$	+1	+1	+1	+1	+1
$A_2$	+1	+1	+1	-1	-1
$B_1$	+1	-1	+1	+1	-1
$B_2$	+1	-1	+1	-1	+1
E	+2	0	-2	0	0



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+1 ← -1 ←	$A_2$	+1	+1	+1	-1	-1
	$B_1$	+1	-1	+1	+1	-1
	$B_2$	+1	-1	+1	-1	+1
	E	+2	0	-2	0	0



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+1 ₊ -1 ₊	$A_2$	+1	+1	+1	-1	-1
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	$B_2$	+1	-1	+1	-1	+1
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	E	+2	0	-2	0	0



#### Character table and irreducible representations (Irrep)









Character table and irreducible representations (Irrep)







**Basis** functions

# **Crystallographic Point Groups**

### [There are 32 crystallographic point groups in 3D]

		Hermann	-Mauguin					
Crystal family	Crystal system	(full)	(short)	Shubnikov	Schoenflies	Orbitoid	Coxeter	Order
Triclinic		1	1	1	<i>C</i> <sub>1</sub>	11	[]+	1
		1	1	2	$C_i = S_2$	×	[2+,2+]	2
		2	2	2	<i>C</i> <sub>2</sub>	22	[2]+	2
Mon	oclinic	m	m	m	$C_s = C_{1h}$	*	[]	2
		$\frac{2}{m}$	2/m	2:m	C <sub>2h</sub>	2*	[2,2+]	4
		222	222	2:2	$D_2 = V$	222	[2,2]+	4
Ortho	rhombic	mm2	mm2	$2 \cdot m$	C <sub>2v</sub>	*22	[2]	4
		$\frac{2}{m}\frac{2}{m}\frac{2}{m}$	mmm	$m \cdot 2:m$	$D_{2h} = V_h$	*222	[2,2]	8
		4	4	4	<i>C</i> <sub>4</sub>	44	[4]+	4
		4	4	Ĩ.	S <sub>4</sub>	2×	[2+,4+]	4
		$\frac{4}{m}$	4/m	4:m	C <sub>4h</sub>	4*	[2,4 <sup>+</sup> ]	8
Tetra	agonal	422	422	4:2	D <sub>4</sub>	422	[4,2]+	8
		4mm	4mm	$4 \cdot m$	C <sub>4v</sub>	*44	[4]	8
		42m	42m	$\tilde{4} \cdot m$	$D_{2d} = V_d$	2*2	[2+,4]	8
		$\frac{4}{m}\frac{2}{m}\frac{2}{m}$	4/mmm	$m \cdot 4:m$	D <sub>4h</sub>	*422	[4,2]	16
	3	3	3	C3	33	[3]+	3	
	Trigonal	3	3	õ	$C_{3i} = S_6$	3×	[2+,6+]	6
		32	32	3:2	D3	322	[3,2]+	6
		3m	3m	$3 \cdot m$	C <sub>3v</sub>	*33	[3]	6
		$\overline{3}\frac{2}{m}$	3m	$ ilde{6} \cdot m$	D <sub>3d</sub>	2*3	[2+,6]	12
		6	6	6	<i>C</i> <sub>6</sub>	66	[6]+	6
Hexagonal		6	6	3:m	C <sub>3h</sub>	3*	[2,3 <sup>+</sup> ]	6
		$\frac{6}{m}$	6/m	6:m	C <sub>6h</sub>	6*	[2,6 <sup>+</sup> ]	12
	Hexagonal	622	622	6:2	D <sub>6</sub>	622	[6,2]+	12
		6mm	6mm	$6 \cdot m$	C <sub>6v</sub>	*66	[6]	12
		6m2	6m2	$m \cdot 3:m$	D <sub>3h</sub>	*322	[3,2]	12
		$\frac{6}{m}\frac{2}{m}\frac{2}{m}$	6/mmm	$m \cdot 6:m$	D <sub>6h</sub>	*622	[6,2]	24
		23	23	3/2	Т	332	[3,3]+	12
		$\frac{2}{m}\overline{3}$	m3	$ ilde{6}/2$	T <sub>h</sub>	3*2	[3+,4]	24
C	ubic	432	432	3/4	0	432	[4,3]+	24
		43m	43m	$3/ ilde{4}$	T <sub>d</sub>	*332	[3,3]	24
		$\frac{4}{m}\overline{3}\frac{2}{m}$	m3m	$ ilde{6}/4$	O <sub>h</sub>	*432	[4,3]	48

C<sub>n</sub>: n-fold rotation C<sub>nh</sub>: C<sub>n</sub> +  $\perp$  mirror C<sub>nv</sub>: C<sub>n</sub> + n || mirrors S<sub>n</sub>: n-fold rotation-reflection D<sub>n</sub>: n-fold rotations + n 2-fold  $\perp$  rotations D<sub>nh</sub>: D<sub>n</sub> +  $\perp$  mirror D<sub>nd</sub>: D<sub>n</sub> + n || mirror T: Tetrahedron [h: with inversion, d: with improper rotations] O: Octahedron [h: with inversion]

#### **Character Tables for Point Groups used in Chemistry**

$C_{3v}$	Е	2 C <sub>3</sub>	3σ <sub>v</sub>
A <sub>1</sub>	1	1	1
A <sub>2</sub>	1	1	-1
E	2	-1	0

#### Symmetry of Rotations and Cartesian products

		Rot         Tr=p         -d          g          i          i          i          i          i          i          i
A <sub>1</sub>	<b>p+d+2f+2g+2h+3i</b> 3j+3k+4l+4m	z, $z^2$ , $x(x^2-3y^2)$ , $z^3$ , $xz(x^2-3y^2)$ , $z^4$ , $xz^2(x^2-3y^2)$ , $z^5$ , $x^2(x^2-3y^2)^2-y^2(3x^2-y^2)^2$ , $xz^3(x^2-3y^2)$ , $z^6$
A <sub>2</sub>	<b>R+f+g+h+2i</b> 2j+2k+3l+3m	$R_{z},  y(3x^{2}-y^{2}),  yz(3x^{2}-y^{2}),  xy(x^{2}-3y^{2})(3x^{2}-y^{2}),  yz^{3}(3x^{2}-y^{2})$
Е	<b>R+p+2d+2f+3g+4h+4i</b> 5j+6k+6l+7m	$\{\mathbf{R}_{x}, \mathbf{R}_{y}\}, \{x, y\}, \{x^{2}-y^{2}, xy\}, \{xz, yz\}, \{z(x^{2}-y^{2}), xyz\}, \{xz^{2}, yz^{2}\}, \{(x^{2}-y^{2})^{2}-4x^{2}y^{2}, xy(x^{2}-y^{2})\}, \{z^{2}(x^{2}-y^{2}), xyz^{2}\}, \{xz^{3}, yz^{3}\}, \{x(x^{2}-(5+2\sqrt{5})y^{2})(x^{2}-(5-2\sqrt{5})y^{2}), y((5+2\sqrt{5})x^{2}-y^{2})((5-2\sqrt{5})x^{2}-y^{2})\}, \{z((x^{2}-y^{2})^{2}-4x^{2}y^{2}), xyz(x^{2}-y^{2})\}, \{z^{3}(x^{2}-y^{2}), xyz^{3}\}, \{xz^{4}, yz^{4}\}, \{xz(x^{2}-(5+2\sqrt{5})y^{2})(x^{2}-(5-2\sqrt{5})y^{2}), yz((5+2\sqrt{5})x^{2}-y^{2})((5-2\sqrt{5})x^{2}-y^{2})\}, \{z^{2}((x^{2}-y^{2})^{2}-4x^{2}y^{2}), xyz^{2}(x^{2}-y^{2})\}, \{z^{4}(x^{2}-y^{2}), xyz^{4}\}, \{xz^{5}, yz^{5}\}$
		Rot         Tr=p         -d        g        h        i

http://gernot-katzers-spice-pages.com/character\_tables/

#### **Character Tables for Point Groups used in Chemistry**

$C_{3v}$	Е	2 C <sub>3</sub>	3σ <sub>v</sub>
A <sub>1</sub>	1	1	1
A <sub>2</sub>	1	1	-1
Е	2	-1	0

### Note: For crystallographic point groups only (32) groups with rotation axes of order n=1,2,3,4,6 are allowed!

#### Symmetry of Rotations and Cartesian products

		Rot         Tr=p        f        g        i        i
A <sub>1</sub>	<b>p+d+2f+2g+2h+3i</b> 3j+3k+4l+4m	z, $z^2$ , $x(x^2-3y^2)$ , $z^3$ , $xz(x^2-3y^2)$ , $z^4$ , $xz^2(x^2-3y^2)$ , $z^5$ , $x^2(x^2-3y^2)^2-y^2(3x^2-y^2)^2$ , $xz^3(x^2-3y^2)$ , $z^6$
A <sub>2</sub>	<b>R+f+g+h+2i</b> 2j+2k+3l+3m	$R_{z}, y(3x^{2}-y^{2}), yz(3x^{2}-y^{2}), yz^{2}(3x^{2}-y^{2}), xy(x^{2}-3y^{2})(3x^{2}-y^{2}), yz^{3}(3x^{2}-y^{2})$
Е	<b>R+p+2d+2f+3g+4h+4i</b> 5j+6k+6l+7m	$\{\mathbf{R}_{x}, \mathbf{R}_{y}\}, \{x, y\}, \{x^{2}-y^{2}, xy\}, \{xz, yz\}, \{z(x^{2}-y^{2}), xyz\}, \{xz^{2}, yz^{2}\}, \{(x^{2}-y^{2})^{2}-4x^{2}y^{2}, xy(x^{2}-y^{2})\}, \{z^{2}(x^{2}-y^{2}), xyz^{2}\}, \{xz^{3}, yz^{3}\}, \{x(x^{2}-(5+2\sqrt{5})y^{2})(x^{2}-(5-2\sqrt{5})x^{2}-y^{2})((5-2\sqrt{5})x^{2}-y^{2})\}, \{z((x^{2}-y^{2})^{2}-4x^{2}y^{2}), xyz(x^{2}-y^{2})\}, \{z^{3}(x^{2}-y^{2}), xyz^{3}\}, \{xz^{4}, yz^{4}\}, \{xz(x^{2}-(5+2\sqrt{5})y^{2})(x^{2}-(5-2\sqrt{5})y^{2}), yz((5+2\sqrt{5})x^{2}-y^{2})((5-2\sqrt{5})x^{2}-y^{2})\}, \{z^{2}((x^{2}-y^{2})^{2}-4x^{2}y^{2}), xyz^{2}(x^{2}-y^{2}), xyz^{3}\}, \{xz^{4}, yz^{4}\}, \{xz(x^{2}-(5+2\sqrt{5})y^{2})(x^{2}-(5-2\sqrt{5})y^{2}), yz((5+2\sqrt{5})x^{2}-y^{2})((5-2\sqrt{5})x^{2}-y^{2})\}, \{z^{2}((x^{2}-y^{2})^{2}-4x^{2}y^{2}), xyz^{2}(x^{2}-y^{2})\}, \{z^{4}(x^{2}-y^{2}), xyz^{4}\}, \{xz^{5}, yz^{5}\}$
	1	Rot         Tr=p          f          i         i

http://gernot-katzers-spice-pages.com/character\_tables/



 $C_n$   $C_1$   $C_2$   $C_3$   $C_4$   $C_5$   $C_6$   $C_7$   $C_8$   $C_9$   $C_{10}$   $C_{11}$   $C_{12}$   $C_{13}$   $C_{14}$   $C_{15}$   $C_{16}$   $C_{17}$   $C_{18}$   $C_{19}$   $C_{20}$  (  $\mathbf{C_{nv}} \quad \mathbf{C_{2v}} \ \mathbf{C_{3v}} \ \mathbf{C_{4v}} \ \mathbf{C_{5v}} \ \mathbf{C_{6v}} \ \mathbf{C_{7v}} \ \mathbf{C_{8v}} \ \mathbf{C_{9v}} \ \mathbf{C_{10v}} \ \mathbf{C_{11v}} \ \mathbf{C_{12v}} \ \mathbf{C_{13v}} \ \mathbf{C_{14v}} \ \mathbf{C_{15v}} \ \mathbf{C_{16v}} \ \mathbf{C_{17v}} \ \mathbf{C_{18v}} \ \mathbf{C_{19v}} \ \mathbf{C_{20v}} \ \mathbf{C_{10v}} \ \mathbf{C_{11v}} \ \mathbf{C_{12v}} \ \mathbf{C_{14v}} \ \mathbf{C_{15v}} \ \mathbf{C_{16v}} \ \mathbf{C_{17v}} \ \mathbf{C_{18v}} \ \mathbf{C_{19v}} \ \mathbf{C_{20v}} \ \mathbf{C_{10v}} \ \mathbf{C_{11v}} \$  $C_{nh} C_s C_{2h} C_{3h} C_{4h} C_{5h} C_{6h} C_{7h} C_{8h} C_{9h} C_{10h} C_{11h} C_{12h} C_{13h} C_{14h} C_{15h} C_{16h} C_{17h} C_{18h} C_{19h} C_{20h} (C_{10h} C_{10h} C_{10h}$  $D_n$   $D_2$   $D_3$   $D_4$   $D_5$   $D_6$   $D_7$   $D_8$   $D_9$   $D_{10}$   $D_{11}$   $D_{12}$   $D_{13}$   $D_{14}$   $D_{15}$   $D_{16}$   $D_{17}$   $D_{18}$   $D_{19}$   $D_{20}$   $D_{10}$   $D_{10}$  DD<sub>nh</sub> D<sub>2h</sub> D<sub>3h</sub> D<sub>4h</sub> D<sub>5h</sub> D<sub>6h</sub> D<sub>7h</sub> D<sub>8h</sub> D<sub>9h</sub> D<sub>10h</sub> D<sub>11h</sub> D<sub>12h</sub> D<sub>13h</sub> D<sub>14h</sub> D<sub>15h</sub> D<sub>16h</sub> D<sub>17h</sub> D<sub>18h</sub> D<sub>19h</sub> D<sub>20h</sub> J D<sub>nd</sub>  $D_{2d} D_{3d} D_{4d} D_{5d} D_{6d} D_{7d} D_{8d} D_{9d} D_{10d} D_{11d} D_{12d} D_{13d} D_{14d} D_{15d} D_{16d} D_{17d} D_{18d} D_{19d} D_{20d} D_{20d} D_{10d} D_{1$ Ci S<sub>6</sub> **S**<sub>16</sub> S<sub>20</sub> S<sub>n</sub> S<sub>4</sub> S<sub>8</sub> **S**<sub>10</sub> **S**<sub>12</sub> **S**<sub>14</sub> S<sub>18</sub> T T<sub>d</sub> T<sub>h</sub> O O<sub>h</sub> isometric I I<sub>h</sub> Schoenflies symbol:



 $C_n$   $C_1$   $C_2$   $C_3$   $C_4$   $C_5$   $C_6$   $C_7$   $C_8$   $C_9$   $C_{10}$   $C_{11}$   $C_{12}$   $C_{13}$   $C_{14}$   $C_{15}$   $C_{16}$   $C_{17}$   $C_{18}$   $C_{19}$   $C_{20}$  (  $\mathbf{C_{nv}} \quad \mathbf{C_{2v}} \ \mathbf{C_{3v}} \ \mathbf{C_{4v}} \ \mathbf{C_{5v}} \ \mathbf{C_{6v}} \ \mathbf{C_{7v}} \ \mathbf{C_{8v}} \ \mathbf{C_{9v}} \ \mathbf{C_{10v}} \ \mathbf{C_{11v}} \ \mathbf{C_{12v}} \ \mathbf{C_{13v}} \ \mathbf{C_{14v}} \ \mathbf{C_{15v}} \ \mathbf{C_{16v}} \ \mathbf{C_{17v}} \ \mathbf{C_{18v}} \ \mathbf{C_{19v}} \ \mathbf{C_{20v}} \ \mathbf{C_{10v}} \ \mathbf{C_{11v}} \ \mathbf{C_{12v}} \ \mathbf{C_{14v}} \ \mathbf{C_{15v}} \ \mathbf{C_{16v}} \ \mathbf{C_{17v}} \ \mathbf{C_{18v}} \ \mathbf{C_{19v}} \ \mathbf{C_{20v}} \ \mathbf{C_{10v}} \ \mathbf{C_{11v}} \$  $C_{nh} C_s C_{2h} C_{3h} C_{4h} C_{5h} C_{6h} C_{7h} C_{8h} C_{9h} C_{10h} C_{11h} C_{12h} C_{13h} C_{14h} C_{15h} C_{16h} C_{17h} C_{18h} C_{19h} C_{20h} (C_{10h} C_{10h} C_{10h}$  $D_n$   $D_2$   $D_3$   $D_4$   $D_5$   $D_6$   $D_7$   $D_8$   $D_9$   $D_{10}$   $D_{11}$   $D_{12}$   $D_{13}$   $D_{14}$   $D_{15}$   $D_{16}$   $D_{17}$   $D_{18}$   $D_{19}$   $D_{20}$   $D_{10}$   $D_{10}$  D $\mathbf{D}_{nh} \qquad \mathbf{D}_{2h} \mathbf{D}_{3h} \mathbf{D}_{4h} \mathbf{D}_{5h} \mathbf{D}_{6h} \mathbf{D}_{7h} \mathbf{D}_{8h} \mathbf{D}_{9h} \mathbf{D}_{10h} \mathbf{D}_{11h} \mathbf{D}_{12h} \mathbf{D}_{13h} \mathbf{D}_{14h} \mathbf{D}_{15h} \mathbf{D}_{16h} \mathbf{D}_{17h} \mathbf{D}_{18h} \mathbf{D}_{19h} \mathbf{D}_{20h} \mathbf{J}_{10h} \mathbf{D}_{11h} \mathbf{D}_{12h} \mathbf{D}_{13h} \mathbf{D}_{14h} \mathbf{D}_{15h} \mathbf{D}_{16h} \mathbf{D}_{17h} \mathbf{D}_{18h} \mathbf{D}_{19h} \mathbf{D}_{20h} \mathbf{J}_{10h} \mathbf{D}_{11h} \mathbf{D}_{12h} \mathbf{D}_{11h} \mathbf{D}_{12h} \mathbf{D}_{11h} \mathbf{D}_{12h} \mathbf{D}_{11h} \mathbf{D}_{12h} \mathbf{D}_{11h} \mathbf{D}_$ D<sub>nd</sub>  $D_{2d} D_{3d} D_{4d} D_{5d} D_{6d} D_{7d} D_{8d} D_{9d} D_{10d} D_{11d} D_{12d} D_{13d} D_{14d} D_{15d} D_{16d} D_{17d} D_{18d} D_{19d} D_{20d} D_{20d} D_{10d} D_{1$ Ci **S**<sub>6</sub> **S**<sub>16</sub> S<sub>20</sub> S<sub>n</sub> S<sub>4</sub> S<sub>8</sub> **S**<sub>10</sub> **S**<sub>12</sub> **S**<sub>14</sub> S<sub>18</sub> T T<sub>d</sub> T<sub>h</sub> O O<sub>h</sub> isometric I I<sub>h</sub> Schoenflies symbol:



 $C_n$   $C_1$   $C_2$   $C_3$   $C_4$   $C_5$   $C_6$   $C_7$   $C_8$   $C_9$   $C_{10}$   $C_{11}$   $C_{12}$   $C_{13}$   $C_{14}$   $C_{15}$   $C_{16}$   $C_{17}$   $C_{18}$   $C_{19}$   $C_{20}$   $C_{20}$  $\mathbf{C_{nv}} \quad \mathbf{C_{2v}} \ \mathbf{C_{3v}} \ \mathbf{C_{4v}} \ \mathbf{C_{5v}} \ \mathbf{C_{6v}} \ \mathbf{C_{7v}} \ \mathbf{C_{8v}} \ \mathbf{C_{9v}} \ \mathbf{C_{10v}} \ \mathbf{C_{11v}} \ \mathbf{C_{12v}} \ \mathbf{C_{13v}} \ \mathbf{C_{14v}} \ \mathbf{C_{15v}} \ \mathbf{C_{16v}} \ \mathbf{C_{17v}} \ \mathbf{C_{18v}} \ \mathbf{C_{19v}} \ \mathbf{C_{20v}} \ \mathbf{C_{10v}} \ \mathbf{C_{11v}} \ \mathbf{C_{12v}} \ \mathbf{C_{14v}} \ \mathbf{C_{15v}} \ \mathbf{C_{16v}} \ \mathbf{C_{17v}} \ \mathbf{C_{18v}} \ \mathbf{C_{19v}} \ \mathbf{C_{20v}} \ \mathbf{C_{10v}} \ \mathbf{C_{11v}} \$  $C_{nh} C_s C_{2h} C_{3h} C_{4h} C_{5h} C_{6h} C_{7h} C_{8h} C_{9h} C_{10h} C_{11h} C_{12h} C_{13h} C_{14h} C_{15h} C_{16h} C_{17h} C_{18h} C_{19h} C_{20h} (C_{10h} C_{10h} C_{10h}$  $D_n$   $D_2$   $D_3$   $D_4$   $D_5$   $D_6$   $D_7$   $D_8$   $D_9$   $D_{10}$   $D_{11}$   $D_{12}$   $D_{13}$   $D_{14}$   $D_{15}$   $D_{16}$   $D_{17}$   $D_{18}$   $D_{19}$   $D_{20}$   $D_{10}$   $D_{10}$  DD<sub>2h</sub> D<sub>3h</sub> D<sub>4h</sub> D<sub>5h</sub> D<sub>6h</sub> D<sub>7h</sub> D<sub>8h</sub> D<sub>9h</sub> D<sub>10h</sub> D<sub>11h</sub> D<sub>12h</sub> D<sub>13h</sub> D<sub>14h</sub> D<sub>15h</sub> D<sub>16h</sub> D<sub>17h</sub> D<sub>18h</sub> D<sub>19h</sub> D<sub>20h</sub> J D<sub>nh</sub> D<sub>nd</sub>  $D_{2d} D_{3d} D_{4d} D_{5d} D_{6d} D_{7d} D_{8d} D_{9d} D_{10d} D_{11d} D_{12d} D_{13d} D_{14d} D_{15d} D_{16d} D_{17d} D_{18d} D_{19d} D_{20d} D_{20d} D_{10d} D_{1$ Ci  $S_6 S_8$ **S**<sub>16</sub> S<sub>20</sub> S<sub>n</sub> S<sub>4</sub> **S**<sub>10</sub> **S**<sub>12</sub> S<sub>14</sub> S<sub>18</sub> T T<sub>d</sub> T<sub>h</sub> O O<sub>h</sub> isometric I I<sub>h</sub> Schoenflies symbol:



 $C_n$   $C_1$   $C_2$   $C_3$   $C_4$   $C_5$   $C_6$   $C_7$   $C_8$   $C_9$   $C_{10}$   $C_{11}$   $C_{12}$   $C_{13}$   $C_{14}$   $C_{15}$   $C_{16}$   $C_{17}$   $C_{18}$   $C_{19}$   $C_{20}$   $C_{20}$  $\mathbf{C_{nv}} \quad \mathbf{C_{2v}} \ \mathbf{C_{3v}} \ \mathbf{C_{4v}} \ \mathbf{C_{5v}} \ \mathbf{C_{6v}} \ \mathbf{C_{7v}} \ \mathbf{C_{8v}} \ \mathbf{C_{9v}} \ \mathbf{C_{10v}} \ \mathbf{C_{11v}} \ \mathbf{C_{12v}} \ \mathbf{C_{13v}} \ \mathbf{C_{14v}} \ \mathbf{C_{15v}} \ \mathbf{C_{16v}} \ \mathbf{C_{17v}} \ \mathbf{C_{18v}} \ \mathbf{C_{19v}} \ \mathbf{C_{20v}} \ \mathbf{C_{10v}} \ \mathbf{C_{11v}} \ \mathbf{C_{12v}} \ \mathbf{C_{14v}} \ \mathbf{C_{15v}} \ \mathbf{C_{16v}} \ \mathbf{C_{17v}} \ \mathbf{C_{18v}} \ \mathbf{C_{19v}} \ \mathbf{C_{20v}} \ \mathbf{C_{10v}} \ \mathbf{C_{11v}} \$  $C_{nh} C_s C_{2h} C_{3h} C_{4h} C_{5h} C_{6h} C_{7h} C_{8h} C_{9h} C_{10h} C_{11h} C_{12h} C_{13h} C_{14h} C_{15h} C_{16h} C_{17h} C_{18h} C_{19h} C_{20h} (C_{10h} C_{10h} C_{10h}$  $D_n$   $D_2$   $D_3$   $D_4$   $D_5$   $D_6$   $D_7$   $D_8$   $D_9$   $D_{10}$   $D_{11}$   $D_{12}$   $D_{13}$   $D_{14}$   $D_{15}$   $D_{16}$   $D_{17}$   $D_{18}$   $D_{19}$   $D_{20}$   $D_{10}$   $D_{10}$  DD<sub>2h</sub> D<sub>3h</sub> D<sub>4h</sub> D<sub>5h</sub> D<sub>6h</sub> D<sub>7h</sub> D<sub>8h</sub> D<sub>9h</sub> D<sub>10h</sub> D<sub>11h</sub> D<sub>12h</sub> D<sub>13h</sub> D<sub>14h</sub> D<sub>15h</sub> D<sub>16h</sub> D<sub>17h</sub> D<sub>18h</sub> D<sub>19h</sub> D<sub>20h</sub> J D<sub>nh</sub> D<sub>nd</sub>  $D_{2d} D_{3d} D_{4d} D_{5d} D_{6d} D_{7d} D_{8d} D_{9d} D_{10d} D_{11d} D_{12d} D_{13d} D_{14d} D_{15d} D_{16d} D_{17d} D_{18d} D_{19d} D_{20d} D_{20d} D_{10d} D_{1$ Ci S<sub>6</sub> S<sub>8</sub> **S**<sub>16</sub>  $S_{20}$ S<sub>n</sub> S<sub>4</sub> **S**<sub>10</sub> **S**<sub>12</sub> S<sub>14</sub> S<sub>18</sub> T T<sub>d</sub> T<sub>h</sub> O O<sub>h</sub> isometric I I<sub>h</sub> Schoenflies symbol:



 $C_n$   $C_1$   $C_2$   $C_3$   $C_4$   $C_5$   $C_6$   $C_7$   $C_8$   $C_9$   $C_{10}$   $C_{11}$   $C_{12}$   $C_{13}$   $C_{14}$   $C_{15}$   $C_{16}$   $C_{17}$   $C_{18}$   $C_{19}$   $C_{20}$  (  $\mathbf{C_{nv}} \quad \mathbf{C_{2v}} \ \mathbf{C_{3v}} \ \mathbf{C_{4v}} \ \mathbf{C_{5v}} \ \mathbf{C_{6v}} \ \mathbf{C_{7v}} \ \mathbf{C_{8v}} \ \mathbf{C_{9v}} \ \mathbf{C_{10v}} \ \mathbf{C_{11v}} \ \mathbf{C_{12v}} \ \mathbf{C_{13v}} \ \mathbf{C_{14v}} \ \mathbf{C_{15v}} \ \mathbf{C_{16v}} \ \mathbf{C_{17v}} \ \mathbf{C_{18v}} \ \mathbf{C_{19v}} \ \mathbf{C_{20v}} \ \mathbf{C_{10v}} \ \mathbf{C_{11v}} \ \mathbf{C_{12v}} \ \mathbf{C_{14v}} \ \mathbf{C_{15v}} \ \mathbf{C_{16v}} \ \mathbf{C_{17v}} \ \mathbf{C_{18v}} \ \mathbf{C_{19v}} \ \mathbf{C_{20v}} \ \mathbf{C_{10v}} \ \mathbf{C_{11v}} \$  $C_{nh} C_s C_{2h} C_{3h} C_{4h} C_{5h} C_{6h} C_{7h} C_{8h} C_{9h} C_{10h} C_{11h} C_{12h} C_{13h} C_{14h} C_{15h} C_{16h} C_{17h} C_{18h} C_{19h} C_{20h} (C_{10h} C_{10h} C_{10h}$  $D_n$   $D_2$   $D_3$   $D_4$   $D_5$   $D_6$   $D_7$   $D_8$   $D_9$   $D_{10}$   $D_{11}$   $D_{12}$   $D_{13}$   $D_{14}$   $D_{15}$   $D_{16}$   $D_{17}$   $D_{18}$   $D_{19}$   $D_{20}$   $D_{10}$   $D_{10}$  D $\mathbf{D}_{nh} \qquad \mathbf{D}_{2h} \mathbf{D}_{3h} \mathbf{D}_{4h} \mathbf{D}_{5h} \mathbf{D}_{6h} \mathbf{D}_{7h} \mathbf{D}_{8h} \mathbf{D}_{9h} \mathbf{D}_{10h} \mathbf{D}_{11h} \mathbf{D}_{12h} \mathbf{D}_{13h} \mathbf{D}_{14h} \mathbf{D}_{15h} \mathbf{D}_{16h} \mathbf{D}_{17h} \mathbf{D}_{18h} \mathbf{D}_{19h} \mathbf{D}_{20h} \mathbf{J}_{10h} \mathbf{D}_{11h} \mathbf{D}_{12h} \mathbf{D}_{13h} \mathbf{D}_{14h} \mathbf{D}_{15h} \mathbf{D}_{16h} \mathbf{D}_{17h} \mathbf{D}_{18h} \mathbf{D}_{19h} \mathbf{D}_{20h} \mathbf{J}_{10h} \mathbf{D}_{11h} \mathbf{D}_{12h} \mathbf{D}_{11h} \mathbf{D}_{12h} \mathbf{D}_{11h} \mathbf{D}_{12h} \mathbf{D}_{11h} \mathbf{D}_{12h} \mathbf{D}_{11h} \mathbf{D}_$ D<sub>nd</sub>  $D_{2d} D_{3d} D_{4d} D_{5d} D_{6d} D_{7d} D_{8d} D_{9d} D_{10d} D_{11d} D_{12d} D_{13d} D_{14d} D_{15d} D_{16d} D_{17d} D_{18d} D_{19d} D_{20d} D_{20d} D_{10d} D_{1$ Ci  $S_6 S_8$ **S**<sub>16</sub>  $S_{20}$ Sn **S**<sub>4</sub> **S**<sub>10</sub> **S**<sub>12</sub> S<sub>14</sub> S<sub>18</sub> T T<sub>d</sub> T<sub>h</sub> O O<sub>h</sub> isometric I I<sub>h</sub> Schoenflies symbol:

# What does this all have to do with SC order parameters?

From fermionic anti-symmetry:

$$\hat{\Delta}(\mathbf{k}) = - \hat{\Delta}^T (-\mathbf{k})$$

 $\Delta_{\alpha\beta}(\mathbf{k}) \sim \left\langle c_{-\mathbf{k}\alpha} c_{\mathbf{k}\beta} \right\rangle$ 

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If inversion is a symmetry:  $P\hat{\Delta}(\mathbf{k})P^{-1} = \hat{\Delta}(-\mathbf{k}) = \pm \Delta(\mathbf{k})$ 

[Assumption: does not modify the internal DOFs]

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Two decoupled sectors of SC order parameters:

$$\hat{\Delta}_{E}(\mathbf{k}) = -\hat{\Delta}_{E}^{T}(-\mathbf{k}) = -\hat{\Delta}_{E}^{T}(\mathbf{k})$$
  
.....  $(i\sigma_{2})$   
 $\sim |\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle$   
Spin Singlet  
Even Parity

 $\hat{\Delta}_{O}(\mathbf{k}) = -\hat{\Delta}_{O}^{T}(-\mathbf{k}) = \hat{\Delta}_{O}^{T}(\mathbf{k})$  $\sigma_3 \propto \sigma_1(i\sigma_2)$   $\sigma_0 \propto \sigma_2(i\sigma_2)$ •  $\sigma_1 \propto \sigma_3(i\sigma_2)$  $\sim |\uparrow\uparrow\rangle - |\downarrow\downarrow\rangle$  $\sim |\uparrow\uparrow\rangle + |\downarrow\downarrow\rangle$  $\sim |\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle$ **Spin triplet Odd Parity** 

From fermionic anti-symmetry:

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For a generic symmetry G:

$$D(G)\hat{\Delta}(\mathbf{k})D(G)^{-1} = \hat{\Delta}[D_{3D}^{-1}(G)\mathbf{k}] = \pm \Delta(\mathbf{k})$$

From fermionic anti-symmetry:

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Can classify the order parameter according to its properties under a given symmetry operation (as even/odd in analogy to the parity)

PreservesBreaksSymmetrySymmetry

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Can classify the order parameter according to its properties under a given symmetry operation (as even/odd in analogy to the parity)



Note: Now there can be multiple symmetry operations present! [Irreducible representations are now useful!]

+1

-1



Basis functions


**Conventional SC: (almost always) Fully gapped!** 

Basis functions

**Special scenario I: 2D Irrep and Nematicity** 

Gap

**1D Irrep** 

...preserves the point group symmetry.

Gap amplitude



...breaks the point group symmetry.

**Special scenario I: 2D Irrep and Nematicity** 

+

**2D** Irrep

... preserves the point group symmetry.

...breaks the point group symmetry.





**Special scenario I: 2D Irrep and Nematicity** 

Gap

**1D** Irrep

**2D Irrep** 

...preserves the point group symmetry.

...breaks the point group symmetry.

Cu<sub>x</sub>Bi<sub>2</sub>Se<sub>3</sub>

[NEMATIC SC]

What are the observable consequences?

+

- Distinct anisotropy in C/T and H<sub>c2</sub>
- Associated lattice deformations



Gap amplitude





### **Special scenario II: 2D Irrep and TRSB**

A complex superposition of the two components in a 2D irrep usually lifts the nodes (generally more stable):



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Note: Isotropic Gap, but certainly unconventional!

What are the observable consequences?

- Polar Kerr Effect
- Muon Spin Relaxation

Sr<sub>2</sub>RuO<sub>4</sub>  $(p_{p_{d}}, 0)$   $(p_{d}, 0)$  $(p_{d}, 0)$ 

J. Xia et al., Phys. Rev. Lett. 97, 167002 (2006)

## $D_{4h} = D_4 + inversion$

$\mathbf{D}_{4\mathbf{h}}_{h=16}$	E	2 C <sub>4</sub>	С <sub>2</sub>	$2C_2^\prime$	$2  C_2^{\prime\prime}$	i	2 S <sub>4</sub>	σ <sub>h</sub>	2 σ <sub>v</sub>	2 σ <sub>d</sub>
A <sub>1g</sub>	1	1	1	1	1	1	1	1	1	1
A <sub>2g</sub>	1	1	1	-1	-1	1	1	1	-1	-1
B <sub>1g</sub>	1	-1	1	1	-1	1	-1	1	1	-1
B <sub>2g</sub>	1	-1	1	-1	1	1	-1	1	-1	1
Eg	2	0	-2	0	0	2	0	-2	0	0
A <sub>1u</sub>	1	1	1	1	1	-1	-1	-1	-1	-1
A <sub>2u</sub>	1	1	1	-1	-1	-1	-1	-1	1	1
B <sub>1u</sub>	1	-1	1	1	-1	-1	1	-1	-1	1
B <sub>2u</sub>	1	-1	1	-1	1	-1	1	-1	1	-1
Eu	2	0	-2	0	0	-2	0	2	0	0



## $D_{4h} = D_4 + inversion$

$\mathbf{D}_{4\mathbf{h}}_{h=16}$	Ε	2 C <sub>4</sub>	С <sub>2</sub>	$2C_2^\prime$	$2  C_2^{\prime\prime}$	i	2 S <sub>4</sub>	σ <sub>h</sub>	2 σ <sub>v</sub>	$2\sigma_{d}$
A <sub>1g</sub>	1	1	1	1	1	1	1	1	1	1
A <sub>2g</sub>	1	1	1	-1	-1	1	1	1	-1	-1
B <sub>1g</sub>	1	-1	1	1	-1	1	-1	1	1	-1
B <sub>2g</sub>	1	-1	1	-1	1	1	-1	1	-1	1
Eg	2	0	-2	0	0	2	0	-2	0	0
A <sub>1u</sub>	1	1	1	1	1	-1	-1	-1	-1	-1
A <sub>2u</sub>	1	1	1	-1	-1	-1	-1	-1	1	1
B <sub>1u</sub>	1	-1	1	1	-1	-1	1	-1	-1	1
B <sub>2u</sub>	1	-1	1	-1	1	-1	1	-1	1	-1
Eu	2	0	-2	0	0	-2	0	2	0	0



#### Symmetry of Rotations and Cartesian products

		Rot         Tr=p         -d        g        i
A <sub>1g</sub>	<b>d+2g+2i</b> 3k+3m	$z^{2}, (x^{2}-y^{2})^{2}-4x^{2}y^{2}, z^{4}, z^{2}((x^{2}-y^{2})^{2}-4x^{2}y^{2}), z^{6}$
A <sub>2g</sub>	<b>R+g+i</b> 2k+2m	$R_z, xy(x^2-y^2), xyz^2(x^2-y^2)$
B <sub>1g</sub>	<b>d+g+2i</b> 2k+3m	$x^{2}-y^{2}, z^{2}(x^{2}-y^{2}), x^{2}(x^{2}-3y^{2})^{2}-y^{2}(3x^{2}-y^{2})^{2}, z^{4}(x^{2}-y^{2})$
B <sub>2g</sub>	<b>d+g+2i</b> 2k+3m	$xy, xyz^2, xy(x^2-3y^2)(3x^2-y^2), xyz^4$
Eg	<b>R+d+2g+3i</b> 4k+5m	$\{\mathbf{R}_{x}, \mathbf{R}_{y}\}, \{xz, yz\}, \{xz(x^{2}-3y^{2}), yz(3x^{2}-y^{2})\}, \{xz^{3}, yz^{3}\}, \{xz(x^{2}-(5+2\sqrt{5})y^{2})(x^{2}-(5-2\sqrt{5})y^{2}), yz((5+2\sqrt{5})x^{2}-y^{2})((5-2\sqrt{5})x^{2}-y^{2})\}, \{xz^{3}(x^{2}-3y^{2}), yz^{3}(3x^{2}-y^{2})\}, \{xz^{5}, yz^{5}\}$
A <sub>1u</sub>	h j+21	
		$xy_{\mathcal{L}}(x - y)$
A <sub>2u</sub>	<b>p+f+2h</b> 2j+31	$z, z^{3}, z((x^{2}-y^{2})^{2}-4x^{2}y^{2}), z^{5}$
A <sub>2u</sub> B <sub>1u</sub>	<b>p+f+2h</b> 2j+31 <b>f+h</b> 2j+21	$z, z^{3}, z((x^{2}-y^{2})^{2}-4x^{2}y^{2}), z^{5}$
A <sub>2u</sub> B <sub>1u</sub> B <sub>2u</sub>	p+f+2h           2j+31           f+h           2j+21           f+h           2j+21	$x_{2}(x^{2}-y^{2}) = 4x^{2}y^{2}, z^{5}$
A <sub>2u</sub> B <sub>1u</sub> B <sub>2u</sub> E <sub>u</sub>	p+f+2h           2j+31           f+h           2j+21           f+h           2j+21           p+2f+3h           4j+51	$x_{1}x_{2}x_{1}y_{2}y_{2}x_{3}z_{3}(x^{2}-y^{2})^{2}-4x^{2}y^{2}, z^{5}$ $x_{2}, z^{3}, z((x^{2}-y^{2})^{2}-4x^{2}y^{2}), z^{5}$ $x_{2}, x_{2}y^{3}$ $x_{2}, x_{2}y^{3}$ $z_{2}(x^{2}-y^{2}), z^{3}(x^{2}-y^{2})$ $z_{3}(x^{2}-y^{2})$ $z_{3}(x^{2}-y^{2}), z^{3}(x^{2}-y^{2})$ $z_{3}(x^{2}-y^{2}), z^{3}(x^{2}-y^{2})$ $z_{3}(x^{2}-y^{2}), z^{3}(x^{2}-y^{2}), z^{3}(x^{2}-y^{2}-y^{2}), z^{3}(x^{2}-y^{2}-y^{2}), z^{3}(x^{2}-y^{2}-$

## $D_{4h} = D_4 + inversion$

<b>D</b> 4h <sub><i>h</i>=16</sub>	Ε	2 C <sub>4</sub>	с <sub>2</sub>	$2C_2^\prime$	$2  C_2^{\prime\prime}$	i	2 S <sub>4</sub>	σ <sub>h</sub>	2 σ <sub>v</sub>	2 σ <sub>d</sub>
A <sub>1g</sub>	1	1	1	1	1	1	1	1	1	1
A <sub>2g</sub>	1	1	1	-1	-1	1	1	1	-1	-1
B <sub>1g</sub>	1	-1	1	1	-1	1	-1	1	1	-1
B <sub>2g</sub>	1	-1	1	-1	1	1	-1	1	-1	1
Eg	2	0	-2	0	0	2	0	-2	0	0
A <sub>1u</sub>	1	1	1	1	1	-1	-1	-1	-1	-1
A <sub>2u</sub>	1	1	1	-1	-1	-1	-1	-1	1	1
B <sub>1u</sub>	1	-1	1	1	-1	-1	1	-1	-1	1
B <sub>2u</sub>	1	-1	1	-1	1	-1	1	-1	1	-1
Eu	2	0	-2	0	0	-2	0	2	0	0



#### Symmetry of Rotations and Cartesian products

Only gives us information about the k-dependent part of the gap function.

A <sub>1g</sub>	<b>d+2g+2i</b> 3k+3m	$z^{2}, (x^{2}-y^{2})^{2}-4x^{2}y^{2}, z^{4}, z^{2}((x^{2}-y^{2})^{2}-4x^{2}y^{2}), z^{6}$
A <sub>2g</sub>	<b>R+g+i</b> 2k+2m	$R_z, xy(x^2-y^2), xyz^2(x^2-y^2)$
B <sub>1g</sub>	<b>d+g+2i</b> 2k+3m	$x^{2}-y^{2}, z^{2}(x^{2}-y^{2}), x^{2}(x^{2}-3y^{2})^{2}-y^{2}(3x^{2}-y^{2})^{2}, z^{4}(x^{2}-y^{2})$
B <sub>2g</sub>	<b>d+g+2i</b> 2k+3m	$xy, xyz^2, xy(x^2-3y^2)(3x^2-y^2), xyz^4$
Eg	<b>R+d+2g+3i</b> 4k+5m	$\{\mathbf{R}_{x}, \mathbf{R}_{y}\}, \{xz, yz\}, \{xz(x^{2}-3y^{2}), yz(3x^{2}-y^{2})\}, \{xz^{3}, yz^{3}\}, \{xz(x^{2}-(5+2\sqrt{5})y^{2})(x^{2}-(5-2\sqrt{5})y^{2}), yz((5+2\sqrt{5})x^{2}-y^{2})((5-2\sqrt{5})x^{2}-y^{2})\}, \{xz^{3}(x^{2}-3y^{2}), yz^{3}(3x^{2}-y^{2})\}, \{xz^{5}, yz^{5}\}$
A <sub>1u</sub>	<b>h</b> j+21	xyz(x <sup>2</sup> -y <sup>2</sup> )
A <sub>2u</sub>	<b>p+f+2h</b> 2j+31	z, $z^3$ , $z((x^2-y^2)^2-4x^2y^2)$ , $z^5$
B <sub>1u</sub>	<b>f+h</b> 2j+21	xyz, xyz <sup>3</sup>
B <sub>2u</sub>	<b>f+h</b> 2j+21	$z(x^2-y^2), z^3(x^2-y^2)$
Eu	<b>p+2f+3h</b> 4j+51	$ \{x, y\}, \{x(x^2-3y^2), y(3x^2-y^2)\}, \{xz^2, yz^2\}, \{x(x^2-(5+2\sqrt{5})y^2)(x^2-(5-2\sqrt{5})y^2), y((5+2\sqrt{5})x^2-y^2)((5-2\sqrt{5})x^2-y^2)\}, \{xz^2(x^2-3y^2), yz^2(3x^2-y^2)\}, \{xz^4, yz^4\} $
		Rot         Tr=p         - d         g         h         i         i

# Phenomenological theory of unconventional superconductivity

Manfred Sigrist and Kazuo Ueda Rev. Mod. Phys. **63**, 239 – Published 1 April 1991

Irreducible representation $\Gamma$	Basis function
$ \begin{array}{c} \Gamma_{1}^{+} \\ \Gamma_{2}^{+} \\ \Gamma_{3}^{+} \\ \Gamma_{4}^{+} \\ \Gamma_{5}^{+} \end{array} $	(a) $\psi(\Gamma_1^+;\mathbf{k}) = 1, k_x^2 + k_y^2, k_z^2$ $\psi(\Gamma_2^+;\mathbf{k}) = k_x k_y (k_x^2 - k_y^2)$ $\psi(\Gamma_3^+;\mathbf{k}) = k_x^2 - k_y^2$ $\psi(\Gamma_4^+;\mathbf{k}) = k_x k_y$ $\psi(\Gamma_5^+,1;\mathbf{k}) = k_x k_z$ $\psi(\Gamma_5^+,2;\mathbf{k}) = k_y k_z$
$\Gamma_1^-$ $\Gamma_2^-$ $\Gamma_3^-$ $\Gamma_4^-$ $\Gamma_5^-$	(b) $\mathbf{d}(\Gamma_{1}^{-};\mathbf{k}) = \mathbf{\hat{x}}k_{x} + \mathbf{\hat{y}}k_{y}, \mathbf{\hat{z}}k_{z}$ $\mathbf{d}(\Gamma_{2}^{-};\mathbf{k}) = \mathbf{\hat{x}}k_{y} - \mathbf{\hat{y}}k_{x}$ $\mathbf{d}(\Gamma_{3}^{-};\mathbf{k}) = \mathbf{\hat{x}}k_{x} - \mathbf{\hat{y}}k_{x}$ $\mathbf{d}(\Gamma_{4}^{-};\mathbf{k}) = \mathbf{\hat{x}}k_{y} + \mathbf{\hat{y}}k_{x}$ $\mathbf{d}(\Gamma_{5}^{-},1;\mathbf{k}) = \mathbf{\hat{x}}k_{z}, \mathbf{\hat{z}}k_{x}$ $\mathbf{d}(\Gamma_{5}^{-},2;\mathbf{k}) = \mathbf{\hat{y}}k_{z}, \mathbf{\hat{z}}k_{y}$

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Ir repre	reducible esentation Γ	Basis function				
	$\Gamma_1^+$ $\Gamma_2^+$	(a) $\psi(\Gamma_1^+;\mathbf{k}) = 1, \ k_x^2 + k_y^2, \ k_z^2$ $\psi(\Gamma_2^+;\mathbf{k}) = k, \ k, \ (k^2 - k^2)$		A <sub>1g</sub>	<b>d+2g+2i</b> 3k+3m	$z^2$ , $(x^2-y^2)^2-4x^2y^2$ , $z^4$ ,
	$\Gamma_{3}^{+}$ $\Gamma_{4}^{+}$ $\Gamma_{4}^{+}$	$\psi(\Gamma_{4}^{+};\mathbf{k}) = k_{x}^{2} - k_{y}^{2}$ $\psi(\Gamma_{4}^{+};\mathbf{k}) = k_{x}k_{y}$ $\psi(\Gamma_{4}^{+};\mathbf{k}) = k_{x}k_{y}$	$\leftrightarrow$	A <sub>2g</sub>	<b>R+g+i</b> 2k+2m	$R_z, xy(x^2-y^2), xyz^2(x^2-y^2)$
	15	$\psi(\Gamma_5, 1, \mathbf{k}) = k_x k_z$ $\psi(\Gamma_5^+, 2; \mathbf{k}) = k_y k_z$		B <sub>1g</sub>	<b>d+g+2i</b> 2k+3m	$x^2-y^2$ , $z^2(x^2-y^2)$ , $x^2(x^2-y^2)$
	$\Gamma_1^-$ $\Gamma_2^-$	(b) $\mathbf{d}(\Gamma_1^-;\mathbf{k}) = \mathbf{\hat{x}}k_x + \mathbf{\hat{y}}k_y, \ \mathbf{\hat{z}}k_z$ $\mathbf{d}(\Gamma_2^-;\mathbf{k}) = \mathbf{\hat{x}}k_y - \mathbf{\hat{y}}k_z$		B <sub>2g</sub>	<b>d+g+2i</b> 2k+3m	$xy, xyz^2, xy(x^2-3y^2)(3x)$
	$\Gamma_3^-$ $\Gamma_4^-$ $\Gamma_5^-$	$ \mathbf{d}(\Gamma_3^-;\mathbf{k}) = \mathbf{\hat{x}}k_x - \mathbf{\hat{y}}k_x  \mathbf{d}(\Gamma_4^-;\mathbf{k}) = \mathbf{\hat{x}}k_y + \mathbf{\hat{y}}k_x  \mathbf{d}(\Gamma_5^-,1;\mathbf{k}) = \mathbf{\hat{x}}k_z, \ \mathbf{\hat{z}}k_x $		Eg	<b>R+d+2g+3i</b> 4k+5m	$\{R_x, R_y\}, \{xz, yz\}, \{z, yz$
		$\mathbf{d}(\Gamma_5, 2; \mathbf{k}) = \mathbf{\hat{y}} k_z, \mathbf{\hat{z}} k_y$				

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Irreducible representation $\Gamma$	Basis function
$ \begin{array}{c} \Gamma_{1}^{+} \\ \Gamma_{2}^{+} \\ \Gamma_{3}^{+} \\ \Gamma_{4}^{+} \\ \Gamma_{5}^{+} \end{array} $	(a) $\psi(\Gamma_1^+;\mathbf{k}) = 1, k_x^2 + k_y^2, k_z^2$ $\psi(\Gamma_2^+;\mathbf{k}) = k_x k_y (k_x^2 - k_y^2)$ $\psi(\Gamma_3^+;\mathbf{k}) = k_x^2 - k_y^2$ $\psi(\Gamma_4^+;\mathbf{k}) = k_x k_y$ $\psi(\Gamma_5^+,1;\mathbf{k}) = k_x k_z$ $\psi(\Gamma_5^+,2;\mathbf{k}) = k_y k_z$
$\Gamma_1^-$ $\Gamma_2^-$ $\Gamma_3^-$ $\Gamma_4^-$ $\Gamma_5^-$	(b) $\mathbf{d}(\Gamma_{1}^{-};\mathbf{k}) = \mathbf{\hat{x}}k_{x} + \mathbf{\hat{y}}k_{y}, \mathbf{\hat{z}}k_{z}$ $\mathbf{d}(\Gamma_{2}^{-};\mathbf{k}) = \mathbf{\hat{x}}k_{y} - \mathbf{\hat{y}}k_{x}$ $\mathbf{d}(\Gamma_{3}^{-};\mathbf{k}) = \mathbf{\hat{x}}k_{x} - \mathbf{\hat{y}}k_{x}$ $\mathbf{d}(\Gamma_{4}^{-};\mathbf{k}) = \mathbf{\hat{x}}k_{y} + \mathbf{\hat{y}}k_{x}$ $\mathbf{d}(\Gamma_{5}^{-},1;\mathbf{k}) = \mathbf{\hat{x}}k_{z}, \mathbf{\hat{z}}k_{x}$ $\mathbf{d}(\Gamma_{5}^{-},2;\mathbf{k}) = \mathbf{\hat{y}}k_{z}, \mathbf{\hat{z}}k_{y}$

Phenomenological theory of unconventional superconductivity

Manfred Sigrist and Kazuo Ueda Rev. Mod. Phys. **63**, 239 – Published 1 April 1991

Irreducible representation Γ	Basis function			
$\Gamma_1^+$ $\Gamma_2^+$	(a) $\psi(\Gamma_1^+;\mathbf{k}) = 1, \ k_x^2 + k_y^2, \ k_z^2$ $\psi(\Gamma_2^+;\mathbf{k}) = k_x k_y (k_x^2 - k_y^2)$	A <sub>1u</sub>	<b>h</b> j+21	$xyz(x^2-y^2)$
$\Gamma_{4}^{+}$ $\Gamma_{5}^{+}$	$\psi(\Gamma_3^+;\mathbf{k}) = k_x^2 - k_y^2$ $\psi(\Gamma_4^+;\mathbf{k}) = k_x k_y$ $\psi(\Gamma_5^+, 1;\mathbf{k}) = k_x k_z$	A <sub>2u</sub>	<b>p+f+2h</b> 2j+31	z, $z^3$ , $z((x^2-y^2)^2-4x^2y^2)$ , z
	$\psi(\Gamma_5^+, 2; \mathbf{k}) = k_y k_z$ (b)	B <sub>1u</sub>	<b>f+h</b> 2j+21	xyz, xyz <sup>3</sup>
$\Gamma_1^-$ $\Gamma_2^-$ $\Gamma_3^-$	$\mathbf{d}(\Gamma_1^-;\mathbf{k}) = \mathbf{\hat{x}}k_x + \mathbf{\hat{y}}k_y, \ \mathbf{\hat{z}}k_z$ $\mathbf{d}(\Gamma_2^-;\mathbf{k}) = \mathbf{\hat{x}}k_y - \mathbf{\hat{y}}k_z$ $\mathbf{d}(\Gamma_3^-;\mathbf{k}) = \mathbf{\hat{x}}k_x - \mathbf{\hat{y}}k_z$	<b>???</b> B <sub>2u</sub>	<b>f+h</b> 2j+21	$z(x^2-y^2), z^3(x^2-y^2)$
$\frac{\Gamma_4^-}{\Gamma_5^-}$	$\mathbf{d}(\Gamma_4^-;\mathbf{k}) = \mathbf{\hat{x}}k_y + \mathbf{\hat{y}}k_x$ $\mathbf{d}(\Gamma_5^-,1;\mathbf{k}) = \mathbf{\hat{x}}k_z, \ \mathbf{\hat{z}}k_x$ $\mathbf{d}(\Gamma_5^-,2;\mathbf{k}) = \mathbf{\hat{x}}k_z, \ \mathbf{\hat{z}}k_x$	Eu	<b>p+2f+3h</b> 4j+51	$\{x, y\}, \{x(x^2-3y^2), y(3x^2-y^2)\}$

Phenomenological theory of unconventional superconductivity

Manfred Sigrist and Kazuo Ueda Rev. Mod. Phys. **63**, 239 – Published 1 April 1991

TABLE IV. (a) Even-parity basis gap functions  $\widehat{\Delta}(\Gamma, m; \mathbf{k}) = i \widehat{\sigma}_y \psi(\Gamma, m; \mathbf{k})$  and (b) odd-parity basis gap functions  $\widehat{\Delta}(\Gamma, m; \mathbf{k}) = i [\widehat{\boldsymbol{\sigma}} \cdot \mathbf{d}(\Gamma, m; \mathbf{k})] \widehat{\sigma}_y$  for the tetragonal lattice symmetry  $(D_{4h})$ .

Irreducible representation Γ	Basis function
$\Gamma_1^+ \\ \Gamma_2^+ \\ \Gamma_3^+ \\ \Gamma_4^+ \\ \Gamma_5^+ \end{cases}$	(a) $\psi(\Gamma_1^+;\mathbf{k}) = 1, k_x^2 + k_y^2, k_z^2$ $\psi(\Gamma_2^+;\mathbf{k}) = k_x k_y (k_x^2 - k_y^2)$ $\psi(\Gamma_3^+;\mathbf{k}) = k_x^2 - k_y^2$ $\psi(\Gamma_4^+;\mathbf{k}) = k_x k_y$ $\psi(\Gamma_5^+, 1;\mathbf{k}) = k_x k_z$ $\psi(\Gamma_5^+, 2;\mathbf{k}) = k_y k_z$
$\Gamma_1^-$ $\Gamma_2^-$ $\Gamma_3^-$ $\Gamma_4^-$ $\Gamma_5^-$	(b) $d(\Gamma_{1}^{-};\mathbf{k}) = \mathbf{\hat{x}}k_{x} + \mathbf{\hat{y}}k_{y}, \mathbf{\hat{z}}k_{z}$ $d(\Gamma_{2}^{-};\mathbf{k}) = \mathbf{\hat{x}}k_{y} - \mathbf{\hat{y}}k_{x}$ $d(\Gamma_{3}^{-};\mathbf{k}) = \mathbf{\hat{x}}k_{x} - \mathbf{\hat{y}}k_{x}$ $d(\Gamma_{4}^{-};\mathbf{k}) = \mathbf{\hat{x}}k_{y} + \mathbf{\hat{y}}k_{x}$ $d(\Gamma_{5}^{-},1;\mathbf{k}) = \mathbf{\hat{x}}k_{z}, \mathbf{\hat{z}}k_{x}$ $d(\Gamma_{5}^{-},2;\mathbf{k}) = \mathbf{\hat{y}}k_{z}, \mathbf{\hat{z}}k_{y}$

### In the presence of SOC:

Symmetry operations also act on the spin DOF and influence the classification of SC order parameters.

Spin singlet (associated with  $\sigma_0$ ) always transforms trivially;

The irreps associated with each spin configuration in the triplet sector can be deduced from the explicit form of the generators:

$$C_{4z} = e^{i\pi\sigma_3/4} = \frac{\sigma_0 - i\sigma_3}{\sqrt{2}}$$
$$C_{2x} = e^{i\pi\sigma_1/2} = i\sigma_1$$

 $P = \sigma_0$  Homework!

# Phenomenological theory of unconventional superconductivity

Manfred Sigrist and Kazuo Ueda

Irreducible

representation  $\Gamma$ 

Rev. Mod. Phys. 63, 239 - Published 1 April 1991

TABLE II. (a) Even-parity basis gap functions  $\hat{\Delta}(\Gamma, m; \mathbf{k}) = i \hat{\sigma}_y \psi(\Gamma, m; \mathbf{k})$  and (b) odd-parity basis gap functions  $\hat{\Delta}(\Gamma, m; \mathbf{k}) = i [\hat{\sigma} \cdot \mathbf{d}(\Gamma, m; \mathbf{k})] \hat{\sigma}_y$  for the cubic lattice symmetry  $(O_h)$ .

(a)

**Basis** functions

TABLE III. (a) Even-parity basis gap functions  $\hat{\Delta}(\Gamma,m;\mathbf{k})=i\hat{\sigma}_y\psi(\Gamma,m;\mathbf{k})$  and (b) odd-parity basis gap functions  $\hat{\Delta}(\Gamma,m;\mathbf{k})=i[\hat{\sigma}\cdot\mathbf{d}(\Gamma,m;\mathbf{k})]\hat{\sigma}_y$  for the hexagonal lattice symmetry  $(D_{6h})$ .

**Basis** functions

Irreducible

representation  $\Gamma$ 

TABLE IV. (a) Even-parity basis gap functions  $\hat{\Delta}(\Gamma, m; \mathbf{k}) = i\hat{\sigma}_y \psi(\Gamma, m; \mathbf{k})$  and (b) odd-parity basis gap functions  $\hat{\Delta}(\Gamma, m; \mathbf{k}) = i[\hat{\boldsymbol{\sigma}} \cdot \mathbf{d}(\Gamma, m; \mathbf{k})]\hat{\sigma}_y$  for the tetragonal lattice symmetry  $(D_{4h})$ .

$\Gamma_1^+$	$\psi(\Gamma_1^+;\mathbf{k})=1, \ k_x^2+k_y^2+k_z^2$	$\Gamma_1^+$
$\Gamma_2^+$	$\psi(\Gamma_2^+;\mathbf{k}) = (k_x^2 - k_y^2)(k_y^2 - k_z^2)(k_z^2 - k_x^2)$	$\Gamma_2^+$ $\Gamma_1^+$
$\Gamma_3^+$	$\psi(\Gamma_3^+, 1; \mathbf{k}) = 2k_z^2 - k_x^2 - k_y^2$ $\psi(\Gamma_3^+, 2; \mathbf{k}) = \sqrt{3}(k^2 - k^2)$	$\Gamma_4^+$
<b>F</b> <sup>+</sup>	$\psi(\Gamma_{3}^{+},2,\mathbf{k}) = k k (k^{2}-k^{2})$	$\Gamma_5^+$
• 4	$\psi(\Gamma_{4}^{+},\mathbf{k}) = k_{z}k_{x}(k_{z}^{2} - k_{x}^{2})$ $\psi(\Gamma_{4}^{+},\mathbf{k}) = k_{z}k_{x}(k_{z}^{2} - k_{x}^{2})$ $\psi(\Gamma_{4}^{+},\mathbf{k}) = k_{x}k_{y}(k_{x}^{2} - k_{y}^{2})$	$\Gamma_6^+$
$\Gamma_5^+$	$\psi(\Gamma_5^+, 1; \mathbf{k}) = k_y k_z$ $\psi(\Gamma_5^+, 2; \mathbf{k}) = k_z k_x$ $\psi(\Gamma_5^+, 3; \mathbf{k}) = k_x k_y$	$\Gamma_1^-$ $\Gamma_2^-$
	(b)	$\Gamma_3^-$
$\Gamma_1^-$	$\mathbf{d}(\boldsymbol{\Gamma}_1^-;\mathbf{k}) = \mathbf{\hat{x}}k_x + \mathbf{\hat{y}}k_y + \mathbf{\hat{z}}k_z$	
$\Gamma_2^-$	$\mathbf{d}(\Gamma_{2}^{-};\mathbf{k}) = \hat{\mathbf{x}}k_{x}(k_{z}^{2} - k_{y}^{2}) + \hat{\mathbf{y}}k_{y}(k_{x}^{2} - k_{z}^{2}) \\ + \hat{\mathbf{z}}k_{z}(k_{y}^{2} - k_{x}^{2})$	$\Gamma_4^-$
$\Gamma_3^-$	$\mathbf{d}(\Gamma_3^-, 1; \mathbf{k}) = 2\mathbf{\hat{z}}k_z - \mathbf{\hat{x}}k_x - \mathbf{\hat{y}}k_y$ $\mathbf{d}(\Gamma_3^-, 2; \mathbf{k}) = \sqrt{3}(\mathbf{\hat{x}}k_x - \mathbf{\hat{y}}k_y)$	$\Gamma_5^-$
$\Gamma_4^-$	$\mathbf{d}(\Gamma_4^-, 1; \mathbf{k}) = \mathbf{\hat{y}}k_z - \mathbf{\hat{z}}k_y$ $\mathbf{d}(\Gamma_4^-, 2; \mathbf{k}) = \mathbf{\hat{z}}k_x - \mathbf{\hat{x}}k_z$	$\Gamma_6^-$
	$\mathbf{d}(1_4, 3; \mathbf{k}) = \mathbf{\hat{k}} k_y - \mathbf{\hat{y}} k_x$	
$\Gamma_5^-$	$d(\Gamma_5^-, 1; \mathbf{k}) = \widehat{\mathbf{y}}k_z + \widehat{\mathbf{z}}k_y$ $d(\Gamma_5^-, 2; \mathbf{k}) = \widehat{\mathbf{z}}k_x + \widehat{\mathbf{x}}k_z$ $d(\Gamma_5^-, 3; \mathbf{k}) = \widehat{\mathbf{x}}k_y + \widehat{\mathbf{y}}k_z$	

#### (a) r+ $\psi(\Gamma_1^+;\mathbf{k})=1, k_x^2+k_y^2, k_z^2$ $\psi(\Gamma_2^+;\mathbf{k}) = k_x k_y (k_x^2 - 3k_y^2)(k_y^2 - 3k_x^2)$ $\psi(\Gamma_3^+;\mathbf{k}) = k_z k_x (k_x^2 - 3k_y^2)$ $\psi(\Gamma_4^+;\mathbf{k}) = k_z k_v (k_v^2 - 3k_x^2)$ $\psi(\Gamma_5^+, 1; \mathbf{k}) = k_x k_z$ $\psi(\Gamma_5^+,2;\mathbf{k})=k_vk_z$ $\psi(\Gamma_6^+, 1; \mathbf{k}) = k_x^2 - k_y^2$ $\psi(\Gamma_6^+, 2; \mathbf{k}) = 2k_x k_y$ $(\mathbf{b})$ $\mathbf{d}(\Gamma_1^-;\mathbf{k}) = \mathbf{\hat{x}}k_x + \mathbf{\hat{y}}k_y, \mathbf{\hat{z}}k_z$ $\mathbf{d}(\Gamma_2^-;\mathbf{k}) = \mathbf{\hat{x}}k_v - \mathbf{\hat{y}}k_x$ $\mathbf{d}(\boldsymbol{\Gamma}_{3}^{-};\mathbf{k}) = \mathbf{\hat{z}}k_{x}(k_{x}^{2}-3k_{y}^{2}),$ $k_z[(k_x^2-k_y^2)\hat{\mathbf{x}}-2k_xk_y\hat{\mathbf{y}}]$ $\mathbf{d}(\boldsymbol{\Gamma}_{4}^{-};\mathbf{k})=\mathbf{\hat{z}}k_{y}(k_{y}^{2}-3k_{x}^{2}),$ $k_z \left[ (k_v^2 - k_x^2) \hat{\mathbf{y}} - 2k_x k_v \hat{\mathbf{x}} \right]$ $\mathbf{d}(\Gamma_5^-, 1; \mathbf{k}) = \mathbf{\hat{x}}k_z, \mathbf{\hat{z}}k_x$ $\mathbf{d}(\Gamma_5^-, 2; \mathbf{k}) = \mathbf{\hat{y}} k_z, \mathbf{\hat{z}} k_y$ $\mathbf{d}(\Gamma_6^-, 1; \mathbf{k}) = \mathbf{\hat{x}} k_x - \mathbf{\hat{y}} k_y$ $d(\Gamma_6^-, 2; \mathbf{k}) = \mathbf{\hat{x}} k_v - \mathbf{\hat{y}} k_x$

Irreducible representation $\Gamma$	Basis function
$\Gamma_1^+ \ \Gamma_2^+ \ \Gamma_3^+ \ \Gamma_5^+$	(a) $\psi(\Gamma_1^+;\mathbf{k}) = 1, \ k_x^2 + k_y^2, \ k_z^2$ $\psi(\Gamma_2^+;\mathbf{k}) = k_x k_y (k_x^2 - k_y^2)$ $\psi(\Gamma_3^+;\mathbf{k}) = k_x^2 - k_y^2$ $\psi(\Gamma_4^+;\mathbf{k}) = k_x k_y$ $\psi(\Gamma_5^+,1;\mathbf{k}) = k_x k_z$ $\psi(\Gamma_5^+,2;\mathbf{k}) = k_y k_z$
$\Gamma_1^-$ $\Gamma_2^-$ $\Gamma_3^-$ $\Gamma_4^-$ $\Gamma_5^-$	(b) $\mathbf{d}(\Gamma_1^-;\mathbf{k}) = \mathbf{\hat{x}}k_x + \mathbf{\hat{y}}k_y, \ \mathbf{\hat{z}}k_z$ $\mathbf{d}(\Gamma_2^-;\mathbf{k}) = \mathbf{\hat{x}}k_y - \mathbf{\hat{y}}k_x$ $\mathbf{d}(\Gamma_3^-;\mathbf{k}) = \mathbf{\hat{x}}k_x - \mathbf{\hat{y}}k_x$ $\mathbf{d}(\Gamma_4^-;\mathbf{k}) = \mathbf{\hat{x}}k_y + \mathbf{\hat{y}}k_x$ $\mathbf{d}(\Gamma_5^-,1;\mathbf{k}) = \mathbf{\hat{x}}k_z, \ \mathbf{\hat{z}}k_z$ $\mathbf{d}(\Gamma_5^-,2;\mathbf{k}) = \mathbf{\hat{y}}k_z, \ \mathbf{\hat{z}}k_y$

#### Can deduce irreps for all other point groups by "symmetry descent"

### Complete classification of SC order parameters from the perspective of point groups!

#### The 32 point groups

No.	Label		Elements
Trici	linic		
1	1	$C_1$	E
2	ī	C,	E, I
Mon	aclinic		
3	2	<i>C</i> .	$E_{\rm e} C_{\rm rel}$
4	m	C. C.	$E, \sigma$
5	2/m	C.,	$E, C_{2}, I, \sigma$
Orth	orhombic	24	
6	222	$D_{2}$	E. Com Com Con
2	mm2	<i>C</i> <sub>1</sub>	$E C_{2x}, \sigma_{2y}, \sigma_{2y}$
8	mmm	D 20	$E_1 C_{2m} C_{2m} C_{2m} I_n \sigma_n \sigma_n$
Tetra	aanal	14	$-1 - 2x_1 + 2y_2 + 2z_2 - 1 - x_3 - y_1 - z_3$
9	4	C.	$E, C_{1}^{\dagger}, C_{2}^{\dagger}, C_{2}^{\dagger}$
10	4	S.	$E_1 S_{421} C_{421} C_{22}$
11	4/m	Can	$E, C_{-}^{+}, C_{-}^{+}, C_{-}^{-}, I, S_{-}, S_{-}^{+}, \sigma$
12	422	DA	$E_{1} C_{2m}^{+}, C_{4m}^{-}, C_{2m}^{-}, C_{2m}^{-}$
13	4mm	CA	$E, C_{4\pi}^+, C_{4\pi}^-, C_{2\pi}, \sigma_{\mu}, \sigma_{\mu}, \sigma_{d\mu}, \sigma_{d\mu}$
14	<b>4</b> 2m	$D_{2d}$	$E_1 S_{42}^+, S_{42}^-, C_{32}^-, C_{23}^-, C_{23}^-, \sigma_{da}^-, \sigma_{db}^-$
15	4/mmm	DAh	$E_1, C_{4-1}^1, C_{4-1}, C_{2-1}, C_{2-1}, C_{2-1}, C_{2-1}, C_{2-1}$
			$I, S_{4z}^-, S_{4z}^+, \sigma_x, \sigma_y, \sigma_y, \sigma_{da}, \sigma_{db}$
Trig	onal		
16	3	С,	$E, C_3^{\dagger}, C_3$
17	3	C 1:	$E_1 C_3^+, C_3^-, I, S_6^-, S_6^+$
18	32	$D_3$	$E, C_3^+, C_3^-, C_{21}, C_{22}^{\prime}, C_{23}^{\prime}$
19	3m	C.3.	$E, C_{3}^{+}, C_{3}^{-}, \sigma_{d1}, \sigma_{d2}, \sigma_{d3}$
20	3m	Dad	$E, C_3^+, C_3^-, C_{21}^+, C_{22}^+, C_{23}^+, I, S_6^-, S_6^+, \sigma_{a1}^-, \sigma_{a2}^-, \sigma_{a3}^-$
Hexa	agonal		
21	6	C.,	$E, C_{6}^{+}, C_{6}^{-}, C_{3}^{+}, C_{3}^{-}, C_{2}^{-}$
22	6	C3h	$E, S_3^-, S_3^+, C_3^-, C_3^-, \sigma_h$
23	6/m	Con	$E, C_6^+, C_6^-, C_3^-, C_3^-, C_2, I, S_3^-, S_3^+, S_6^-, S_6^+, \sigma_h$
24	622	$D_6$	$E, C_6^1, C_6, C_3^1, C_3, C_2, C_{21}^{\prime}, C_{22}^{\prime}, C_{23}^{\prime}, C_{21}^{\prime}, C_{22}^{\prime}, C_{23}^{\prime}, C_{21}^{\prime}, C_{22}^{\prime}, C_{23}^{\prime}$
25	6mm	$C_{6v}$	$E, C_6^+, C_6^-, C_2^-, C_3^-, C_2^-, \sigma_{d1}^-, \sigma_{d2}^-, \sigma_{d3}^-, \sigma_{v2}^-, \sigma_{v3}^-$
26	<u>6</u> 2m	$D_{3k}$	$E, S_3^-, S_3^+, C_3^+, C_3^-, \sigma_k, C_{21}', C_{22}', C_{23}', \sigma_{v1}, \sigma_{v2}, \sigma_{v3}$
27	6/mmm	Don	$E, C_6^+, C_6^-, C_3^+, C_3^-, C_2, C_{21}', C_{22}', C_{23}', C_{21}', C_{22}', C_{23}'$
			$I, S_3, S_3^+, S_6^-, S_6^+, \sigma_h, \sigma_{d1}, \sigma_{d2}, \sigma_{d3}, \sigma_{v1}, \sigma_{v2}, \sigma_{v3}$
Cubi	ic .		
28	23	Т	$E_{1}, C_{2m}, C_{3j}, C_{3j}$
29	m3	$T_{h}$	$E, C_{2m}, C_{3j}^{-}, C_{3j}^{-}, I, \sigma_m, S_{6j}^{-}, S_{6j}^{+}$
30	432	0	$E, C_{2m}, C_{3j}^+, C_{3j}^-, C_{2p}^-, C_{4m}^+, C_{4m}^-$
31	43m	T <sub>d</sub>	$E, C_{2m}, C_{3j}^+, C_{3j}^-, \sigma_{dp}, S_{4m}^-, S_{4m}^+$
32	m3m	O <sub>h</sub>	$E, C_{2m}, C_{3j}, C_{3j}, C_{2p}, C_{4m}^{+}, C_{4m}^{-}$
			$I, \sigma_m, S_{6j}^-, S_{6j}^+, \sigma_{dp}, S_{4m}^-, S_{4m}^-$



## Have we covered everything? Is the Sigrist-Ueda classification "complete"?

"Yes and No!"



## Have we covered everything? Is the Sigrist-Ueda classification "complete"?

[Generalizations]

I) Multiple internal DOF

**II)** Nonsymmorphic systems [Space group]

## Have we covered everything? Is the Sigrist-Ueda classification "complete"?

[Generalizations]

I) Multiple internal DOF

**II)** Nonsymmorphic systems [Space group]

How to describe the superconducting states in complex materials with multiple internal DOFs?

## **Considering multiple internal DOF (orbitals/sublattice)**

Annica's Lecture:

The mean-field BdG Hamiltonian

$$\hat{H}_{BdG}(\mathbf{k}) = \begin{pmatrix} \hat{H}_0(\mathbf{k}) & \hat{\Delta}(\mathbf{k}) \\ \hat{\Delta}^{\dagger}(\mathbf{k}) & -\hat{H}_0^*(-\mathbf{k}) \end{pmatrix}$$

$$\hat{\Delta}(\mathbf{k}) = \sum_{ab} d_{ab}(\mathbf{k})\hat{\tau}_a \otimes \hat{\sigma}_b(i\hat{\sigma}_2)$$
Orbital/SL Spin

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Orbital/SL Spin

In principle parametrised in terms of (3+1)x(3+1) = 16 functions d<sub>ab</sub>(k)

 $egin{aligned} \sigma_1 &= \sigma_{\mathrm{x}} = egin{pmatrix} 0 & 1 \ 1 & 0 \end{pmatrix} \ \sigma_2 &= \sigma_{\mathrm{y}} = egin{pmatrix} 0 & -i \ i & 0 \end{pmatrix} \ \sigma_3 &= \sigma_{\mathrm{z}} = egin{pmatrix} 1 & 0 \ 0 & -1 \end{pmatrix} \end{aligned}$ 

If a = 0,3: Intra-orbital/SL If a = 1,2: Inter-orbital/SL

If b = 0: Spin Singlet If b = 1,2,3: Spin Triplet

$$\hat{\Delta}(\mathbf{k}) = \sum_{ab} d_{ab}(\mathbf{k})\hat{\tau}_a \otimes \hat{\sigma}_b(i\hat{\sigma}_2)$$

[a,b]	$\hat{ au}_a$	$\hat{\sigma}_b(i\sigma_2)$	Matrix	k
[0,0]	S	Α	А	Е
[0,1]	S	S	S	0
[0,2]	S	S	S	0
[0,3]	S	S	S	0
[1,0]	S	А	А	Е
[1,1]	S	S	S	0
[1,2]	S	S	S	0
[1,3]	S	S	S	0
[2,0]	A	Α	S	0
[2,1]	A	S	Α	Е
[2,2]	A	S	Α	Е
[2,3]	A	S	Α	Е
[3,0]	S	Α	Α	Е
[3,1]	S	S	S	0
[3,2]	S	S	S	0
[3,3]	S	S	S	0

Annica's Lecture:

$$\hat{\Delta}(\mathbf{k}) = -\hat{\Delta}^T(-\mathbf{k})$$

If the matrix is anti-symmetric: k-even If the matrix is symmetric: k-odd

$$\hat{\Delta}(\mathbf{k}) = \sum_{ab} d_{ab}(\mathbf{k})\hat{\tau}_a \otimes \hat{\sigma}_b(i\hat{\sigma}_2)$$

<u>-</u>		전 15년 14일 5년 14일 5년		
[a,b]	$\hat{ au}_a$	$\hat{\sigma}_b(i\sigma_2)$	Matrix	k
[0,0]	S	Α	Α	Е
[0,1]	S	S	S	0
[0,2]	S	S	S	0
[0,3]	S	S	S	0
[1,0]	S	Α	Α	Е
[1,1]	S	S	S	0
[1,2]	S	S	S	0
[1,3]	S	S	S	0
[2,0]	A	Α	S	0
[2,1]	A	S	Α	Е
[2,2]	A	S	Α	Е
[2,3]	A	S	Α	Е
[3,0]	S	Α	Α	Е
[3,1]	S	S	S	0
[3,2]	S	S	S	0
[3,3]	S	S	S	0

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Inversion symmetry:

Equal parity: 
$$P = \pm \hat{\tau}_0 \otimes \hat{\sigma}_0$$

**Opposite parity:**  $P = \hat{\tau}_3 \otimes \hat{\sigma}_0$ 

[a,b]	$\hat{ au}_a$	$\hat{\sigma}_b(i\sigma_2)$	Matrix	k
[0,0]	S	A	Α	E
[0,1]	S	S	S	0
[0,2]	S	S	S	0
[0,3]	S	S	S	0
[1,0]	S	A	А	Е
[1,1]	S	S	S	0
[1,2]	S	S	S	0
[1,3]	S	S	S	0
[2,0]	A	A	S	0
[2,1]	A	S	А	Е
[2,2]	A	S	Α	Е
[2,3]	A	S	Α	Е
[3,0]	S	Α	Α	Е
[3,1]	S	S	S	0
[3,2]	S	S	S	0
[3,3]	S	S	S	0

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				]	
[a,b]	$\hat{ au}_a$	$\hat{\sigma}_b(i\sigma_2)$	Matrix	k	EP
[0, 0]	S	Α	A	Е	Е
[0,1]	S	S	S	0	0
[0,2]	S	S	S	0	0
[0,3]	S	S	S	0	0
[1, 0]	S	A	A	E	Е
[1, 1]	S	S	S	0	0
[1,2]	S	S	S	0	0
[1,3]	S	S	S	0	0
[2, 0]	A	А	S	0	0
[2, 1]	A	S	Α	Е	Е
[2,2]	A	S	Α	Е	Е
[2, 3]	A	S	Α	E	Е
[3,0]	S	Α	Α	Е	Е
[3,1]	S	S	S	0	0
[3,2]	S	S	S	0	0
[3, 3]	S	S	S	0	0

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[a,b]	$\hat{ au}_a$	$\hat{\sigma}_b(i\sigma_2)$	Matrix	k	EP	OP
[0,0]	S	Α	Α	Е	Е	Е
[0,1]	S	S	S	0	0	0
[0,2]	S	S	S	0	0	0
[0,3]	S	S	S	0	0	0
[1, 0]	S	А	Α	Е	Е	0
[1,1]	S	S	S	0	0	Е
[1,2]	S	S	S	0	0	Е
[1,3]	S	S	S	0	0	Е
[2, 0]	A	Α	S	0	0	Е
[2, 1]	A	S	Α	Е	Е	0
[2,2]	A	S	Α	Е	Е	0
[2, 3]	A	S	Α	Е	Е	0
[3,0]	S	Α	Α	Е	Е	Е
[3,1]	S	S	S	0	0	0
[3,2]	S	S	S	0	0	0
[3,3]	S	S	S	0	0	0

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[a,b]	$\hat{ au}_a$	$\hat{\sigma}_b(i\sigma_2)$	Matrix	k	EP	OP	$\mathbf{SL}$
[0, 0]	S	A	Α	Е	Е	Е	Е
[0,1]	S	S	S	0	0	0	0
[0,2]	S	S	S	0	0	0	0
[0,3]	S	S	S	0	0	0	0
[1,0]	S	Α	Α	Е	Е	0	Е
[1,1]	S	S	S	0	0	Е	0
[1,2]	S	S	S	0	0	Е	0
[1,3]	S	S	S	0	0	Е	0
[2, 0]	A	A	S	0	0	Е	Е
[2,1]	A	S	Α	Е	Е	0	0
[2,2]	A	S	Α	Е	Е	0	0
[2,3]	A	S	Α	Е	Е	0	0
[3,0]	S	A	Α	Е	Е	Е	0
[3,1]	S	S	S	0	0	0	Е
[3,2]	S	S	S	0	0	0	Е
[3,3]	S	S	S	0	0	0	Е

$$\hat{\Delta}(\mathbf{k}) = \sum_{ab} d_{ab}(\mathbf{k})\hat{\tau}_a \otimes \hat{\sigma}_b(i\hat{\sigma}_2)$$

[a,b]	$\hat{ au}_a$	$\hat{\sigma}_b(i\sigma_2)$	Matrix	k	EP	OP	$\mathbf{SL}$
[0, 0]	S	A	Α	Е	Е	Е	Е
[0,1]	S	S	S	0	0	0	0
[0,2]	S	S	S	0	0	0	0
[0,3]	S	S	S	0	0	0	0
[1, 0]	S	A	Α	Е	Е	0	Е
[1, 1]	S	S	S	0	0	Е	0
[1, 2]	S	S	S	0	0	Е	0
[1, 3]	S	S	S	0	0	Е	0
[2, 0]	A	A	S	0	0	Е	Е
[2, 1]	A	S	Α	Е	Е	0	0
[2, 2]	A	S	Α	Е	Е	0	0
[2, 3]	A	S	Α	Е	Е	0	0
[3, 0]	S	A	A	Е	Е	Е	0
$\overline{[3,1]}$	S	S	S	0	0	0	Е
[3,2]	S	S	S	0	0	0	Е
[3,3]	S	S	S	0	0	0	Е

$$\hat{\Delta}(\mathbf{k}) = \sum_{ab} d_{ab}(\mathbf{k})\hat{\tau}_a \otimes \hat{\sigma}_b(i\hat{\sigma}_2)$$

### Spin Singlet [b=0]

[a,b]	$\hat{ au}_a$	$\hat{\sigma}_b(i\sigma_2)$	Matrix	k	EP	OP	$\mathbf{SL}$
[0, 0]	S	А	A	Е	Е	Е	Е
[0,1]	S	S	S	0	0	0	0
[0,2]	S	S	S	0	0	0	0
[0,3]	S	S	S	0	0	0	0
[1,0]	S	А	А	Е	Е	0	Е
[1,1]	S	S	S	0	0	Е	0
[1,2]	S	S	S	0	0	Е	0
[1,3]	S	S	S	0	0	Е	0
[2,0]	A	А	S	0	0	Е	Е
[2,1]	A	S	Α	Е	Е	0	0
[2,2]	A	S	Α	Е	Е	0	0
[2,3]	A	S	Α	Е	Е	0	0
[3,0]	S	Α	А	Е	Е	Е	0
[3,1]	S	S	S	0	0	0	Е
[3,2]	S	S	S	0	0	0	Е
[3,3]	S	S	S	0	0	0	Е

$$\hat{\Delta}(\mathbf{k}) = \sum_{ab} d_{ab}(\mathbf{k})\hat{\tau}_a \otimes \hat{\sigma}_b(i\hat{\sigma}_2)$$

Spin Singlet [b=0]

Spin Triplet [b=1,2,3]

	[a,b]	$\hat{ au}_a$	$\hat{\sigma}_b(i\sigma_2)$	Matrix	k	EP	OP	$\mathbf{SL}$
	[0, 0]	S	A	A	Е	Е	Е	Е
	[0, 1]	S	S	S	0	0	0	0
	[0, 2]	S	S	S	0	0	0	0
	[0, 3]	S	S	S	0	0	0	0
	[1, 0]	S	A	А	Е	Е	0	Е
	[1, 1]	S	S	S	0	0	Е	0
	[1, 2]	S	S	S	0	0	Е	0
	[1, 3]	S	S	S	0	0	Е	0
$\rightarrow$	[2, 0]	A	A	S	0	0	Е	Е
$\rightarrow$	[2, 1]	A	S	А	Е	Е	0	0
$\rightarrow$	[2, 2]	A	S	А	Е	Е	0	0
	[2, 3]	A	S	А	Е	Е	0	0
	[3, 0]	S	A	А	Е	Е	Е	0
	[3, 1]	S	S	S	0	0	0	Е
	[3,2]	S	S	S	0	0	0	Е
	[3, 3]	S	S	S	0	0	0	Е

$$\hat{\Delta}(\mathbf{k}) = \sum_{ab} d_{ab}(\mathbf{k})\hat{\tau}_a \otimes \hat{\sigma}_b(i\hat{\sigma}_2)$$

Spin Singlet [b=0]

Spin Triplet [b=1,2,3]

k-dependence does not uniquely define the parity of the SC order parameter!

	[a,b]	$\hat{ au}_a$	$\hat{\sigma}_b(i\sigma_2)$	Matrix	k	EP	OP	$\mathbf{SL}$
	[0,0]	S	А	А	Е	Е	Е	Е
	[0,1]	S	S	S	0	0	0	0
	[0,2]	S	S	S	0	0	0	0
	[0,3]	S	S	S	0	0	0	0
	[1, 0]	S	А	А	Е	Е	0	Е
	[1, 1]	S	S	S	0	0	Е	0
	[1, 2]	S	S	S	0	0	Е	0
	[1, 3]	S	S	S	0	0	Е	0
$\rightarrow$	[2, 0]	A	А	S	0	0	Е	Е
	[2, 1]	A	S	А	Е	Е	0	0
	[2, 2]	A	S	А	Е	Е	0	0
	[2, 3]	A	S	А	Е	Е	0	0
	[3,0]	S	A	А	Е	Е	Е	0
	[3,1]	S	S	S	0	0	0	Е
	[3,2]	S	S	S	0	0	0	Е
	[3, 3]	S	S	S	0	0	0	Е
## Some examples of nontrivial phenomenology

$$\hat{\Delta}(\mathbf{k}) = -\hat{\Delta}^{T}(-\mathbf{k}) \xrightarrow{\text{Only spin}} \hat{\Delta}(\mathbf{k}) = d_{a}(\mathbf{k})\hat{\sigma}_{a}(i\hat{\sigma}_{2})$$

$$\xrightarrow{\hat{\Delta}(\mathbf{k}) = d_{ab}(\mathbf{k})\hat{\tau}_{a} \otimes \hat{\sigma}_{b}(i\hat{\sigma}_{2})$$

$$\hat{\Delta}(\mathbf{k}) = d_{ab}(\mathbf{k})\hat{\tau}_{a} \otimes \hat{\sigma}_{b}(i\hat{\sigma}_{2})$$
Orbital/Layer/Sublattice+Spin

Can transform non-trivially under inversion!

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Can transform non-trivially under inversion!

#### The case of CeRh<sub>2</sub>As<sub>2</sub>

**Sublattice structure** 

$$P = \hat{\tau}_1$$



 $\hat{\Delta}(\mathbf{k}) = d_{33}(\mathbf{k})\hat{\tau}_3 \otimes \hat{\sigma}_3(i\hat{\sigma}_2)$ 

Even-parity, k-odd, intra-layer, spin-triplet

#### Two superconducting phases!

D. Möckli and A. Ramires, Phys. Rev. Research 3, 023204 (2021)



The case of CeRh<sub>2</sub>As<sub>2</sub>

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Even-parity, k-odd, intra-layer, spin-triplet

#### Two superconducting phases!

The case of d-Bi<sub>2</sub>Se<sub>3</sub>

### **Even- and odd-P orbitals**

$$P = \hat{\tau}_3$$



$$\hat{\Delta}(\mathbf{k}) = d_0 \hat{\tau}_1 \otimes \hat{\sigma}_0(i\hat{\sigma}_2)$$

Odd-parity, s-wave, inter-orbital, spin-singlet

#### **Generalized Anderson's Theorem**



The case of CeRh<sub>2</sub>As<sub>2</sub>

**Sublattice structure** 

 $P = \hat{\tau}_1$ 



 $\hat{\Delta}(\mathbf{k}) = d_{33}(\mathbf{k})\hat{\tau}_3 \otimes \hat{\sigma}_3(i\hat{\sigma}_2)$ 

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Odd-parity, s-wave, inter-orbital, spin-singlet

#### **Generalized Anderson's Theorem**

The case of Sr<sub>2</sub>RuO<sub>4</sub> 3 orbitals



Chiral d-wave superconductivity [Orbital antisymmetric spin-triplet]

### Chiral d-wave in 2D FS!

3 t<sub>2g</sub> orbitals/3 bands system



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S. Beck, A. Ramires et al., Phys. Rev. Research 4, 023060 (2022)

3 t<sub>2g</sub> orbitals/3 bands system



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 $3 t_{2g}$  orbitals/3 bands system



## SC states [Even-parity sector]

Irrep	[a,b]	Orbital	Spin
	[0, 0]	symmetric	singlet
4.	[8, 0]	symmetric	singlet
$\Lambda_{1g}$	[4, 3]	antisymmetric	triplet
	[5,2] - [6,1]	antisymmetric	triplet
$A_{2g}$	[5,1] + [6,2]	antisymmetric	triplet
B.	[7, 0]	symmetric	singlet
$D_{1g}$	[5,2] + [6,1]	antisymmetric	triplet
Ba	[1, 0]	symmetric	singlet
$D_{2g}$	[5,1] - [6,2]	antisymmetric	triplet
	$\{[3,0],-[2,0]\}$	symmetric	singlet
$E_g$	$\{[4,2],-[4,1]\}$	antisymmetric	triplet
	$\{[5,3],[6,3]\}$	antisymmetric	triplet

Microscopic basis: E-parity/S-Triplet Band basis: pseudospin-S

 $3 t_{2g}$  orbitals/3 bands system



## SC states [Even-parity sector]

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	$A_{2g}$	[5,1] + [6,2]	antisymmetric	triplet
	$B_{1g}$	[7, 0]	symmetric	singlet
		[5,2] + [6,1]	antisymmetric	triplet
	P.	[1, 0]	symmetric	singlet
	$D_{2g}$	[5,1] - [6,2]	antisymmetric	triplet
		$\{[3,0],-[2,0]\}$	symmetric	singlet
	$E_g$	$\{[4,2],-[4,1]\}$	antisymmetric	triplet
		$\{[5,3],[6,3]\}$	antisymmetric	triplet

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Hund's interaction [inter-orbital]

 $3 t_{2g}$  orbitals/3 bands system



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	$A_{2g}$	[5,1] + [6,2]	antisymmetric	triplet
	$B_{1g}$	[7, 0]	symmetric	singlet
		[5,2] + [6,1]	antisymmetric	triplet
	P.	[1, 0]	symmetric	singlet
	$D_{2g}$	[5,1] - [6,2]	antisymmetric	triplet
		$\{[3,0],-[2,0]\}$	symmetric	singlet
	$E_g$	$\{[4,2],-[4,1]\}$	antisymmetric	triplet
		$\{[5,3],[6,3]\}$	antisymmetric	triplet

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	$A_{1g}$	[4, 3]	antisymmetric	triplet
		[5,2] - [6,1]	antisymmetric	triplet
	$A_{2g}$	[5,1] + [6,2]	antisymmetric	triplet
	$B_{1g}$	[7, 0]	symmetric	singlet
		[5,2] + [6,1]	antisymmetric	triplet
	R <sub>a</sub>	[1, 0]	symmetric	singlet
	$D_{2g}$	[5,1] - [6,2]	antisymmetric	triplet
	$E_g$	$\{[3,0],-[2,0]\}$	symmetric	singlet
		$\{[4,2],-[4,1]\}$	antisymmetric	triplet
		$\{[5,3],[6,3]\}$	antisymmetric	triplet

Microscopic basis: E-parity/S-Triplet Band basis: pseudospin-S





### Hund's interaction [inter-orbital]

- Uncovered mechanism for chiral d-wave!
- Engineering the normal state to enhance T<sub>c</sub>!

S. Beck, A. Ramires et al., Phys. Rev. Research 4, 023060 (2022)

Pz-like orbitals in a quintuple layer



## Pz-like orbitals in a quintuple layer

K-independent sector



Irrep	Spin	Orbital	Parity	Matrix Form
A	Singlet	Trivial	Even	$\hat{ au}_0\otimes\hat{\sigma}_0(i\hat{\sigma}_2)$
	Singlet	IIIviai	Even	$\hat{ au}_3\otimes\hat{\sigma}_0(i\hat{\sigma}_2)$
$A_{1u}$	Triplet	Singlet	Odd	$\hat{ au}_2\otimes\hat{\sigma}_3(i\hat{\sigma}_2)$
$A_{2u}$	Singlet	Triplet	Odd	$\hat{ au}_1\otimes\hat{\sigma}_0(i\hat{\sigma}_2)$
E	Triplet	Singlet	Odd	$i\hat{ au}_2\otimes\hat{\sigma}_1(i\hat{\sigma}_2)$
$E_u$	Tiblet		Ouu	$\hat{ au}_2\otimes\hat{\sigma}_2(i\hat{\sigma}_2)$

Odd parity  $\Rightarrow$  Nodes! [Sensitive to disorder]

### Pz-like orbitals in a quintuple layer

K-independent sector

Irrep	Spin	Orbital	Parity	Matrix Form
<i>A</i> 1	Singlet	Trivial	Even	$\hat{ au}_0 \otimes \hat{\sigma}_0(i\hat{\sigma}_2)$
$T_{1g}$	Singlet			$\hat{ au}_3 \otimes \hat{\sigma}_0(i\hat{\sigma}_2)$
$A_{1u}$	Triplet	Singlet	Odd	$\hat{ au}_2\otimes\hat{\sigma}_3(i\hat{\sigma}_2)$
$A_{2u}$	Singlet	Triplet	Odd	$\hat{ au}_1 \otimes \hat{\sigma}_0(i\hat{\sigma}_2)$
E	Triplet	Singlet	Odd	$i\hat{ au}_2\otimes\hat{\sigma}_1(i\hat{\sigma}_2)$
$E_u$	Inplet	Singlet	Ouu	$\hat{ au}_2 \otimes \hat{\sigma}_2(i\hat{\sigma}_2)$

Odd parity  $\Rightarrow$  Nodes! [Sensitive to disorder]

#### Experiment/Theory



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M. P. Smylie et al., PRB 96, 115145 (2017)
```



## Pz-like orbitals in a quintuple layer

Cu<sub>x</sub>/Nb<sub>x</sub>/Sr<sub>x</sub>(PbSe)<sub>x</sub> T T P1z+ P2z-Se(2) Bi(1') Se(1') Se(1')

#### K-independent sector

Irrep	Spin	Orbital	Parity	Matrix Form
<i>A</i> 1	Singlet	nglot Trivial		$\hat{ au}_0\otimes\hat{\sigma}_0(i\hat{\sigma}_2)$
$T_{1g}$			Liven	$\hat{ au}_3\otimes\hat{\sigma}_0(i\hat{\sigma}_2)$
$A_{1u}$	Triplet	Singlet	Odd	$\hat{ au}_2\otimes\hat{\sigma}_3(i\hat{\sigma}_2)$
$A_{2u}$	Singlet	Triplet	Odd	$\hat{ au}_1\otimes\hat{\sigma}_0(i\hat{\sigma}_2)$
$F_{\rm c}$	Triplet	Singlet	Odd	$i\hat{ au}_2\otimes\hat{\sigma}_1(i\hat{\sigma}_2)$
$L_{u}$	Tublet	Singlet	Ouu	$\hat{ au}_2\otimes\hat{\sigma}_2(i\hat{\sigma}_2)$

Odd parity  $\Rightarrow$  Nodes! [Sensitive to disorder]

### Experiment/Theory









## "Generalised Anderson's Theorem"

L. Andersen\*, A. Ramires\* et al., Sci. Adv. 6, eaay6502 (2020) B. Zinkl and A. Ramires, Phys. Rev. B 106, 014515 (2022)

## Elena's Lecture

# Two sublattices/layers: CeRh<sub>2</sub>As<sub>2</sub>

## **Cartoon picture:**



"Trivial" Even-parity SC



B<sub>z</sub>-robust Odd-parity SC

M. Sigrist et al., J. Phys. Soc. Jpn. 83, 061014 (2014)
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D. Maruyama et al., J. Phys. Soc. Jpn. 81, 034702 (2012)
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Bz



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Khim et al., Science 373, 1012 (2021)

## Successfully addresses the magnetic field anisotropy



Landaeta et al., PRX 12, 031001 (2022)

## Some common themes...



S. Khim et al., Science **373**, 1012 (2021) S. Adenwalla et a., Phys. Rev. Lett. **65**, 2298 (1990) D. Aoki et al., J. Phys. Soc. Jpn. **89**, 053705 (2020) S. Ran et al., Nature Physics 15, 1250 (2019)

## Some common themes...



- Phase diagrams with multiple SC phases are rare!
- Only observed in other two HF materials!
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## Some common themes...



- Phase diagrams with multiple SC phases are rare!
- Only observed in other two HF materials!
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- Common theme: sublattice DOF?!

T. Hazra et al., Phys. Rev. Lett. 130, 136002 (2023)

S. Khim et al., Science **373**, 1012 (2021) S. Adenwalla et a., Phys. Rev. Lett. **65**, 2298 (1990) D. Aoki et al., J. Phys. Soc. Jpn. **89**, 053705 (2020) S. Ran et al., Nature Physics 15, 1250 (2019)



## Bibliography [group theory & superconductivity]

## Phenomenological Theory of Unconventional Superconductivity

Manfred Sigrist and Kazuo Ueda Rev. Mod. Phys **63**, 239 (1991)

## Symmetry aspects of Chiral Superconductors

Aline Ramires Contemporary Physics **63**(2), 71 (2022)

## Nonunitary Superconductivity in Complex Quantum Materials

Aline Ramires J. Phys.: Condens. Matter **34** 304001(2022)

## Still mystery after all these years -- Unconventional SC of Sr<sub>2</sub>RuO<sub>4</sub>

Yoshiteru Maeno, Shingo Yonezawa, Aline Ramires arXiv:2402.12117 [Invited review to appear in JPSJ]

## Have we covered everything? Is the Sigrist-Ueda classification "complete"?

[Generalizations]

I) Multiple internal DOF

**II)** Nonsymmorphic systems [Space group]

A general space-group operation can be written as [Seitz notation]:

 $\{G|\mathbf{t}\}\$ Point

operation

Translation

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Translation

## **Examples:**

$\{E \mid 0\}$	Identity
$\{G   0\}$	Pure point operation
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**Composition:** 

 $\{G_1 | t_1\} \{G_2 | t_2\} = \{G_1 \cdot G_2 | G_1 t_2 + t_1\}$ 

## [There are 230 space groups in 3D]

Ne	Crystal system,		Point group				Space groups (international short
<u>Ne</u>	(count), Bravais lattice	Int'l	Schön.	Orbifold	Cox.	Ord.	symbol)
1	Triclinic (2)	1	C <sub>1</sub>	11	[]+	1	P1
2	$\alpha \alpha \beta^{c}$	1	C <sub>i</sub>	1×	[2+,2+]	2	Pī
3–5	Monoclinic	2	C <sub>2</sub>	22	[2]+	2	P2, P2 <sub>1</sub> C2
6–9	(13)	m	Cs	*11	[]	2	Pm, Pc Cm, Cc
10–15	and be and be a second	2/m	C <sub>2h</sub>	2*	[2,2+]	4	P2/m, P2 <sub>1</sub> /m C2/m, P2/c, P2 <sub>1</sub> /c C2/c
16–24		222	D <sub>2</sub>	222	[2,2]+	4	P222, P222 <sub>1</sub> , P2 <sub>1</sub> 2 <sub>1</sub> 2, P2 <sub>1</sub> 2 <sub>1</sub> 2 <sub>1</sub> , C222 <sub>1</sub> , C222, F222, I222, I2 <sub>1</sub> 2 <sub>1</sub> 2 <sub>1</sub> 2 <sub>1</sub>
25–46	Orthorhombic (59)	mm2	C <sub>2v</sub>	*22	[2]	4	Pmm2, Pmc2 <sub>1</sub> , Pcc2, Pma2, Pca2 <sub>1</sub> , Pnc2, Pmn2 <sub>1</sub> , Pba2, Pna2 <sub>1</sub> , Pnn2 Cmm2, Cmc2 <sub>1</sub> , Ccc2, Amm2, Aem2, Ama2, Aea2 Fmm2, Fdd2 Imm2, Iba2, Ima2
47–74		mmm	D <sub>2h</sub>	*222	[2,2]	8	Pmmm, Pnnn, Pccm, Pban, Pmma, Pnna, Pmna, Pcca, Pbam, Pccn, Pbcm, Pnnm, Pmmn, Pbcn, Pbca, Pnma Cmcm, Cmce, Cmmm, Cccm, Cmme, Ccce Fmmm, Fddd Immm, Ibam, Ibca, Imma

Continues with tetragonal, trigonal, hexagonal and cubic...

https://en.wikipedia.org/wiki/Space\_group

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Note: Order of the SG is infinite (translations!)

Definition: A glide plane consists of a reflection followed by a (non-primitive) translation parallel to the plane of reflection.

1D Example

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Elena's Lecture

3D Example: Elemental Te [helical chains]

3D Example: CeRh<sub>2</sub>As<sub>2</sub>



#### [P3<sub>1</sub>21]



#### Jairo's Lecture

## Altermagnetism: Example of RuO<sub>2</sub>



 $[I, C_{2z}, ...]$ 

[P4<sub>2</sub>/mnm] Nonsymmorphic

https://phys.org/news/2022-11-anomalous-hall-effect-altermagnetic-ruthenium.html

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# Nonsymmorphic Space groups

Consider a space group G with operations  $\{G | t\}$  which leave a given lattice invariant. We can rewrite each operation as:

 $\{G | \mathbf{t}\} = \{G | \mathbf{T}_{PLV} + \tilde{\mathbf{t}}\} = \{E | \mathbf{T}_{PLV}\}\{G | \tilde{\mathbf{t}}\}\$ 

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If, by ANY choice of origin we find that AT LEAST ONE of the elements of G have  $\tilde{\mathbf{t}} \neq 0$  the space group is called NONSYMMORPHIC

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 $\Rightarrow$  Need to take translations into account!

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Translations in 3D: three sets of (infinite) Abelian subgroups

 $\Rightarrow$  Infinite conjugacy classes = Infinite irreps

 $\Rightarrow$  Bloch's functions are basis functions for the group of translations (labelled by momenta **k**)

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Symmorphic groups:  $D_k^{\Gamma_i}(\{R_\alpha | \mathbf{R}_n\}) = e^{i\mathbf{k} \cdot \mathbf{R}_n} D^{\Gamma_i}(R_\alpha),$  $\chi_{k}^{\Gamma_{i}}(\{R_{\alpha}|\boldsymbol{R}_{n}\}) = \mathrm{e}^{\mathrm{i}\boldsymbol{k}\cdot\boldsymbol{R}_{n}}\chi^{\Gamma_{i}}(R_{\alpha}).$ 

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Nonsymmorphic groups: More complicated...but there are tables!



## Nonsymmorphic symmetry Manifestation #1: Symmetry-protected band crossings

s-states in a diamond lattice [Fd-3m]



Hourglass fermions in KHgSb [P6<sub>3</sub>/mmc]



S. M. Young et al., PRL 108, 140405 (2012)

S. Wang et al., Nature 532, 189 (2016)

### Nonsymmorphic symmetry Manifestation #2: New OP connectivities in modulated systems

### **Band perspective**

Fe-based SC [P4/nmm]



Cvektovik et al., Phys. Rev. B 88, 134510 (2013)

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## **OP** perspective

CeRh<sub>2</sub>As<sub>2</sub> [P4/nmm]

If a multi-component order parameter:

 $F_c = \gamma M_1 M_2 P + \dots$ 

$$E_{1/2m} \otimes E_{1/2m} = A_{1g} \oplus A_{2u} \oplus B_{2g} \oplus B_{1u}$$
$$E_{3/4m} \otimes E_{3/4m} = A_{1g} \oplus A_{1u} \oplus B_{2u} \oplus B_{2g}$$



A. Ramires and A. Szabo, arXiv.2309.05664 (2013)

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"there are no line nodes in odd-parity superconductors in the presence of SOC"

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M. R. Norman, PRB 52, 15093 (1995)

## Nonsymmorphic symmetry Manifestation #3: New nodes at the BZ edge



M. R. Norman, PRB 52, 15093 (1995)

S. Kobayashi et al., PRB 94, 134512 (2016)
T. Micklitz et al., PRL 118, 207001 (2017)
T. Micklitz et al., PRB 95, 024508 (2017)
S. Sumita Ph.D. Thesis (2019)

## Summary/Conclusion

#### **Brief introduction to group theory concepts:**

Group  $\Rightarrow$  Conjugacy Classes  $\Rightarrow$  Group Representation  $\Rightarrow$  Character  $\Rightarrow$  Irreducible Representations

#### **Crystallographic Point Groups:**

- $\Rightarrow$  SC order parameter classification
- $\Rightarrow$  Conventional/unconventional
- $\Rightarrow$  Nematic/Chiral

#### **Beyond the Sigrist-Ueda Classification:**

- ⇒ Multiple internal DOFs (orbitals/layers/sublattices)
- $\Rightarrow$  Nonsymmorphic symmetries

"Loopholes" to what we have thought were very well-established concepts and theorems in the field...are there more of them?

Homework!